

FROM EULER, RITZ AND GALERKIN TO MODERN COMPUTING

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Abstract. The so-called Ritz-Galerkin method is one of the most fundamental tools of modern computing. Its origins lie in Hilbert’s “direct” approach to the variational calculus of Euler-Lagrange and in the thesis of Walther Ritz, who died 100 years ago at the age of 31 after a long struggle against tuberculosis. The thesis was submitted in 1902 in Göttingen, in a period of dramatic developments in Physics. Ritz tried to explain the phenomenon of Balmer series in spectroscopy using eigenvalue problems of partial differential equations on rectangular domains. While this physical model quickly turned out to be completely obsolete, his mathematics enabled him later to solve difficult problems in applied sciences. He thereby revolutionized the variational calculus and became one of the fathers of modern computational mathematics. We will see in this article that the path leading to modern computational methods and theory was a long struggle over three centuries requiring the efforts of many great mathematicians.

1. The Variational Calculus of Euler and Lagrange. The most well-known contribution of Walther Ritz is the development of a systematic approach for solving *variational problems*. His methods transformed variational calculus from a theoretical tool to one of practical importance, and are the precursor to many algorithms for modern scientific computing. Going back in history, we will see how variational calculus started in 1696 with the famous challenge concerning the Brachystochrone problem, which led to endless disputes between the Bernoulli brothers. Euler [12] gave in 1744 a general solution to variational problems in the form of a differential equation. Eleven years later, the nineteen-year-old Lagrange then communicated, in a famous letter to Euler, an elegant justification for this equation. The prodigious contributions of Euler concerning the analytic and numerical solutions of differential equations, in particular his *Institutiones Calculi Integralis* [13] from 1668–1670, then added the finishing touch to the theory.

1.1. The Brachystochrone Problem. In 1696, Johann Bernoulli challenged the mathematical world (which included his brother Jacob) with the following problem (see Fig. 1.1): *given two fixed points A and B in a vertical plane, find a curve AMB such that a body gliding on it under gravity, starting from A, arrives after the shortest possible time at B, or in beautiful Latin: Datis in plano verticali duobus punctis A & B, assignare Mobili M viam AMB, per quam gravitate sua descendens, & moveri incipiens a puncto A, brevissimo tempore perveniat ad alterum punctum B.* Already Galileo, in 1638, knew that the *shortest* path, the straight line between A and B, is not the *quickest* and he guessed wrongly that a circle would be the best solution. So the true shape of the best curve remained a mystery for nearly a century.

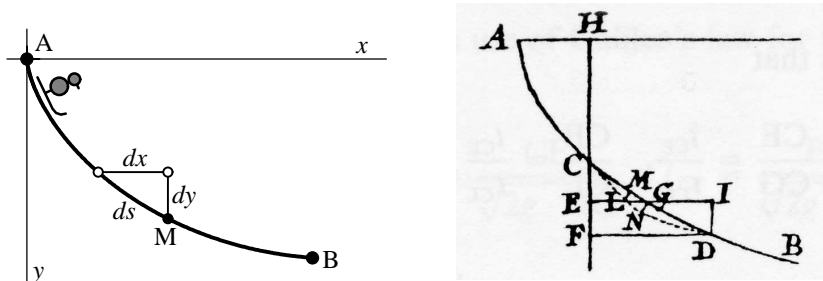


FIG. 1.1. The Brachystochrone problem (left); Jacob Bernoulli’s solution (right)

In order to turn this physical question into a mathematical one¹, we notice that the time for covering a small arc length ds is $dJ = \frac{ds}{v}$. Inserting $v = \sqrt{2gy}$ (Galileo), we obtain the problem of finding a function $y(x)$ with $y(a) = A$, $y(b) = B$ such that the integral

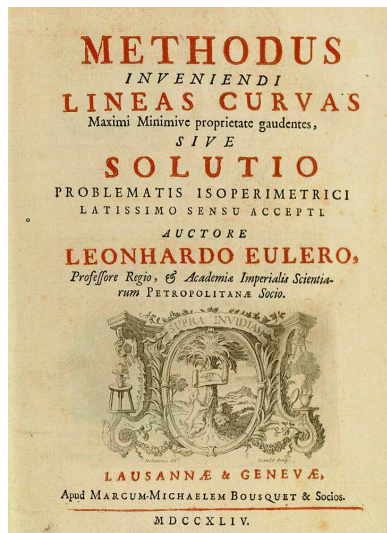
$$(1.1) \quad J = \int_a^b \frac{\sqrt{dx^2 + dy^2}}{\sqrt{2gy}} = \int_a^b \frac{\sqrt{1+p^2}}{\sqrt{2gy}} dx \longrightarrow \min \quad \text{where } p = \frac{dy}{dx}.$$

We will see below that the optimal solution is a cycloide, which allowed Johann Bernoulli to proudly correct an error of Galilei.

1.2. Euler's Differential Equation for a Variational Problem. Half a century later appeared Euler's famous *Methodus inveniendi lineas curvas* [12] (E65, 1744) which solved a great variety of minimization problems, among which

$$(1.2) \quad J = \int_a^b Z(x, y, p) dx \longrightarrow \min/\max \quad \text{where } p = \frac{dy}{dx}$$

where Z is an arbitrary function; see Fig. 1.2 ².



Variational Calc. 1744 E65



Zoom: above envy !

FIG. 1.2. Euler's legacy for the theory of variational calculus, with a zoom.

Euler's Theorem. For an optimal solution we have

$$(1.3) \quad \boxed{N - \frac{d}{dx}P = 0} \quad \text{where } N = \frac{\partial Z}{\partial y}, \quad P = \frac{\partial Z}{\partial p}.$$

¹Letter of de l'Hôpital to Joh. Bernoulli, June 15th, 1696: "Ce probleme me paroist des plus curieux et des plus jolis que l'on ait encore proposé et je serois bien aise de m'y appliquer ; mais pour cela il seroit necessaire que vous me l'envoyassiez reduit à la mathématique pure, car le phisque m'embarasse ...".

²To represent his mathematical colleagues as dogs barking at a theory (the brachystochrone), nailed out of their reach high up on a tree, was not Euler's style, but it was Johann Bernoulli's, who had produced the picture for his own *Opera omnia* in 1742. The editor, M. Bousquet, found it a good idea to reuse the printing plate two years later. Perhaps Bousquet thought that Euler was an ape sitting still higher up in the tree and eating up all the coconuts.

We will come back later to Euler's proof of this formula. Because the term P is in general a function of $p = y'$, and because this term is differentiated once again in Euler's equation, we usually obtain a *second order* ordinary differential equation for $y(x)$, which may be difficult to solve. In the case where Z is *independent of x* , i.e., if $dZ = N dy + P dp$, Euler found, in §30 of caput II of E65, an elegant way to reduce the order: multiply

$$N - \frac{dP}{dx} = 0 \quad \text{by} \quad dy = p dx \quad \Rightarrow \quad N dy - p dP = 0,$$

add and subtract $P dp$ to obtain

$$(1.4) \quad \underbrace{N dy + P dp}_{dZ} - \underbrace{(p dP + P dp)}_{d(p \cdot P)} = 0 \quad \Rightarrow \quad \boxed{Z - p \cdot \frac{\partial Z}{\partial p} = \text{Const.}}$$

Example 1 (smallest curve length, smallest energy; Euler E65, Caput II, §33). We search for a curve with $y(a) = A$, $y(b) = B$ of shortest arc length, i.e.,

$$(1.5) \quad J = \int_a^b \sqrt{1 + p^2} dx \rightarrow \min$$

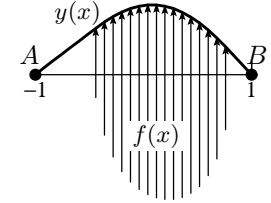
Here, $N = 0$ and we obtain from (1.3) that $\frac{dP}{dx} = 0$, i.e., $P = \frac{p}{\sqrt{1+p^2}}$ must be constant, thus $p = \text{constant}$ too. Not surprisingly, the solution is a straight line. Euler explains that the same result is valid for *any* problem where Z depends uniquely on p . An interesting case is

$$(1.6) \quad J = \int_a^b \frac{p^2}{2} dx \rightarrow \min \quad \text{for which (1.3) becomes} \quad -\frac{d^2y}{dx^2} = 0.$$

This represents an approximation of J in (1.5) for small p and models a stretched flexible cord. If a transversal force $f(x)$ acts on the cord, we get the problem

$$(1.7) \quad J = \int_a^b \left(\frac{p^2}{2} - f \cdot y \right) dx \rightarrow \min$$

which leads to

$$-\frac{d^2y}{dx^2} = f(x).$$


This case was too simple for Euler to be mentioned, but its extension to higher dimensions will become very important later.

Example 2 (the Brachistochrone problem; Euler E65, Caput II, §34). We obtain from (1.4)

$$(1.8) \quad \frac{\sqrt{1 + p^2}}{\sqrt{2gy}} - \frac{p^2}{\sqrt{2gy} \sqrt{1 + p^2}} = C \quad \text{or} \quad 1 = \sqrt{1 + p^2} \sqrt{2gy} \cdot C.$$

It is still not a trivial matter to find a curve with this property. We remark that Johann Bernoulli, with one of his typically brilliant intuitions, arrived immediately at this last equation by applying Snell's law of light refraction. Because $\frac{\sin \alpha}{v} = \text{const} = \frac{1}{v} \frac{dx}{ds} = 1/(\sqrt{1 + p^2} \sqrt{2gy})$, by (1.8) this law is satisfied everywhere and represents, by Fermat's principle, the quickest path.

1.3. Euler's Integral Calculus. A considerable part of Euler's tremendous work is devoted to analytical and numerical methods for the solution of integrals and differential equations. This work culminates in the three volumes of *Institutiones Calculi Integralis* [13] (E342, E366, E385, published 1768-1770). Let us apply this methods to find the solution of the Brachystochrone problem (1.8): we solve the equation for $p = \frac{dy}{dx}$ and obtain

$$(1.9) \quad 1 + p^2 = \frac{1}{C^2 y} \quad \text{or} \quad p = \frac{dy}{dx} = \sqrt{\frac{1}{C^2 y} - 1} = \sqrt{\frac{c-y}{y}}$$

with $c = C^{-2}$. In § 397 of E342, Euler explains that the variables x and y should be separated, i.e.

$$\sqrt{\frac{y}{c-y}} \cdot dy = dx \quad \text{such that} \quad \int \sqrt{\frac{y}{c-y}} \cdot dy = x.$$

The easiest method for such integrals are trigonometric substitutions, which Euler mentioned relatively late in E342 (in § 329). If we set $y = c \cdot \sin^2 u$ and profit from $1 - \sin^2 u = \cos^2 u$, then the integral becomes an easy one and leads to the solution

$$x = cu - \frac{c}{2} \sin 2u \quad y = c \cdot \sin^2 u = \frac{c}{2} - \frac{c}{2} \cos 2u,$$

so that “curvam quaesitam esse *Cycloidem*”.

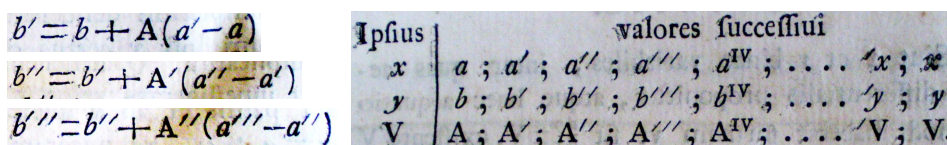


FIG. 1.3. First publication of Euler's numerical method for differential equations.

Numerical Solution. Whenever such an analytical solution is not possible, Euler proposes (in E342, § 650) to compute the solution *vero proxime assignare*, by writing an equation such as (1.9) in the general form

$$(1.10) \quad \frac{dy}{dx} = V(x, y) \quad \text{so that} \quad x_{i+1} = x_i + \Delta x, \quad y_{i+1} = y_i + \Delta x \cdot V(x_i, y_i)$$

are, for $i = 1, 2, 3, \dots$, *valores successivi* of numerical approximations to the solution. The smaller we choose the “step size” Δx , the better will be the numerical approximation. We see in Fig. 1.3 the first publication of this method, and in Fig. 1.4 (left) the numerical result by computing 16 steps with this formula for the Brachystochrone problem³. This is the first of the so-called *finite difference methods*, which dominated all scientific calculations for two centuries.

1.4. Joseph Louis de Lagrange. On August 12th, 1755, the nineteen-year-old Ludovico de la Grange Tournier wrote to *Vir amplissime atque celeberrime L. Euler*, from whom *Vir praestantissime atque excellentissime Lagrange* received a kind and enthusiastic answer (6 sept. 1755).

³The formula does not work for $i = 0$, because V becomes infinite there. Such singularities must be treated separately.

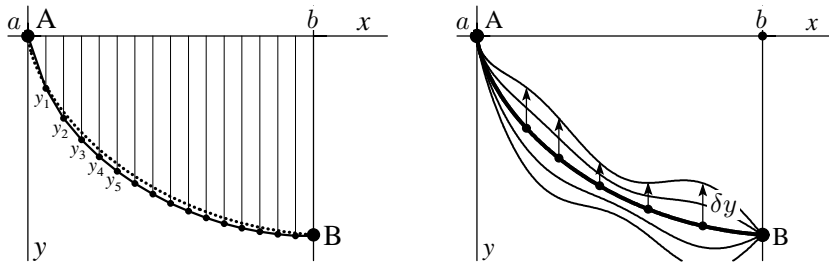


FIG. 1.4. Numerical solution of the Brachystochrone by Euler's method (exact solution dotted, left); Lagrange's variation $\epsilon \delta y$ of y for $\epsilon = -1, -\frac{1}{2}, \frac{1}{2}, 1$ (right).

Lagrange's idea was the following: we perturb the optimal solution $y(x)$ by an arbitrary *variation*⁴ δy . But only another 15 years later became it clear to Euler [14] how "facilis" was the geometric meaning behind Lagrange's formalism: we add a fixed function $\delta(x)$ multiplied by ϵ to $y(x)$ (see Fig. 1.4, right) and insert the result into (1.2). This integral must be minimal for all functions $\delta y(x)$ at $\epsilon = 0$, i.e. the derivative of

$$(1.11) \quad J(\epsilon) = \int_a^b Z(x, y + \epsilon \delta y, p + \epsilon \delta p) dx \rightarrow \min$$

with respect to ϵ must be zero at $\epsilon = 0$. We differentiate:

$$(1.12) \quad \frac{\partial J(\epsilon)}{\partial \epsilon} \Big|_{\epsilon=0} = \int_a^b (N \cdot \delta y + P \cdot \delta p) dx = 0.$$

Since δp is the derivative of δy , we can integrate by parts:

$$(1.13) \quad \int_a^b \left(N - \frac{d}{dx} P \right) \cdot \delta y \cdot dx = 0. \quad \text{"weak solution"}$$

Because $\delta y(a) = \delta y(b) = 0$ (the end points of y are fixed), the end-point contributions of (1.13) vanish. Since δy is arbitrary, we conclude that for all x

$$(1.14) \quad N - \frac{d}{dx} P = 0 \quad \text{"strong solution"}.$$

This last step, a trivial conclusion for the young discoverer, later caused the greatest difficulties.

The passage from the variational problem (1.11) by differentiation to (1.12), then by partial integration to (1.13), finally to (1.14) is the *central highway of variational calculus*. In modern times (1.12) will be called the "directional derivative" of J and the function $\delta y(x)$ in (1.13) a "test function"; this equation will be the starting point of the "Galerkin method". If we then manage to solve equation (1.14), we will have solved the original variational problem *à la Euler*. A very important application of the theory appeared in 1788 with Lagrange's *Mécanique analytique*, where the above "highway" connects the *Lagrangian* of a mechanical system (difference of potential and kinetic energy) to the differential equations of its motion. This later led to Hamiltonian mechanics.

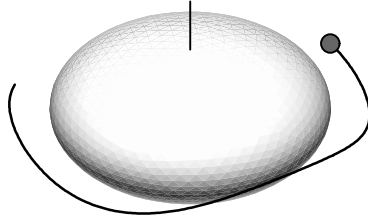


FIG. 2.1. *Attraction of a non-spherical celestial body*

2. Laplace's Equation and Dirichlet's Principle. Since Newton's *Principia* from 1687, we know that celestial bodies move under the action of forces obeying the inverse square law, and since Euler's work (in particular E112 from 1747) their motion obeys second order differential equations with the corresponding term $f_x \approx \frac{x-\xi}{r^3}$, $r = \sqrt{(x-\xi)^2 + (y-\eta)^2 + (z-\zeta)^2}$. But this is only valid for pointwise, or at most perfectly spherical, bodies. If our body has another shape (see Fig. 2.1), the attractive forces become

$$(2.1) \quad f_x = \iiint \rho(\xi, \eta, \zeta) \frac{x-\xi}{r^3} d\xi d\eta d\zeta,$$

and similarly for f_y and f_z , which are cumbersome expressions. Here Laplace (*Théorie des attractions des sphéroïdes et de la figure de la terre*, 1785; see also *Oeuvres I, Mécanique céleste*, p. 157) had the idea to introduce the *Potential* function

$$(2.2) \quad V = \iiint \rho(\xi, \eta, \zeta) \frac{1}{r} d\xi d\eta d\zeta, \quad \text{so that} \quad f_x = -\frac{\partial V}{\partial x},$$

and similarly for f_y and f_z . If we differentiate (2.1) once again with respect to x (and y and z respectively) we find for V the elegant expression

$$(2.3) \quad \Delta V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0,$$

which bears the name *Laplace's equation*.

Soon after, this equation found many other important applications, not only in heaven, but also down on earth:

- theory of *stationary heat transfer* (Fourier 1822);
- theory of *magnetism* (Gauss and Weber in Göttingen, C.F. Gauss [16], Werke 5, p. 195, 1839);
- theory of *electric fields* (W. Thomson, later Lord Kelvin [29], Liouville J. 12, 1847);
- *conformal mappings* (Gauss 1825, Werke IV, p. 189);
- in *complex analysis* (Cauchy 1825, Riemann [36] Thesis 1851);
- irrotational *fluid motion* in 2D (Helmholtz 1858).

All these applications led young Riemann to the following enthusiastic statement:

⁴From here originates the name of the whole theory.

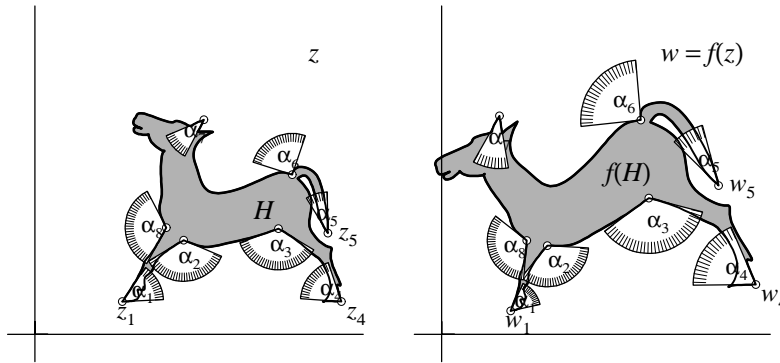


FIG. 2.2. A conformal mapping

“... eine vollkommen in sich abgeschlossene mathematische Theorie ... , welche ... fortschreitet, ohne zu scheiden, ob es sich um die Schwerkraft, oder die Electricität, oder den Magnetismus, oder das Gleichgewicht der Wärme handelt.” (Manuscript of Riemann, Nov. 1850, *Werke* p. 545.)

2.1. Conformal Maps and the Riemann Mapping Theorem. With his thesis [36] from 1851, Riemann founded *geometric function theory*, a subject which studies theorems in complex analysis through elegant geometric considerations, and which was later perfected mainly through the work of Richard Courant (cf. the second part of Hurwitz and Courant [23]). The starting points are the Cauchy-Riemann differential equations for continuously differentiable functions $f(z) = u(x, y) + iv(x, y)$ with $z = x + iy$,

$$(2.4) \quad \frac{\partial v}{\partial y} = \frac{\partial u}{\partial x} \quad \frac{\partial u}{\partial y} = -\frac{\partial v}{\partial x}$$

in Riemann’s hand-writing⁵, which, when differentiated, give

$$(2.5) \quad \Delta u = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0 \quad \text{and} \quad \Delta v = \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} = 0,$$

i.e., u and v are *harmonic*. Furthermore, equations (2.4) tell us that the Jacobian in \mathbb{R}^2 of a continuously differentiable function is of the form

$$(2.6) \quad \begin{pmatrix} \frac{\partial u}{\partial x} & \frac{\partial u}{\partial y} \\ \frac{\partial v}{\partial x} & \frac{\partial v}{\partial y} \end{pmatrix} = \begin{pmatrix} a & b \\ -b & a \end{pmatrix} = \text{Const} \cdot \begin{pmatrix} \cos \phi & \sin \phi \\ -\sin \phi & \cos \phi \end{pmatrix}$$

i.e., it consists locally of a uniform scaling and an orthogonal rotation. Such everywhere angle-preserving functions are called *conformal* (“in kleinsten Theilen ähnlich”; see Fig. 2.2).

Riemann mapping theorem. *For any simply connected domain Ω there exists a bijective conformal mapping to the unit disk $f : \Omega \rightarrow B$. The point z_0 mapped to the origin can be freely chosen in the interior of Ω , as well as the image of one point*

⁵courtesy of Prof. E. Neuenschwander, *Riemanns Einführung in die Funktionentheorie*, Göttingen (1996)

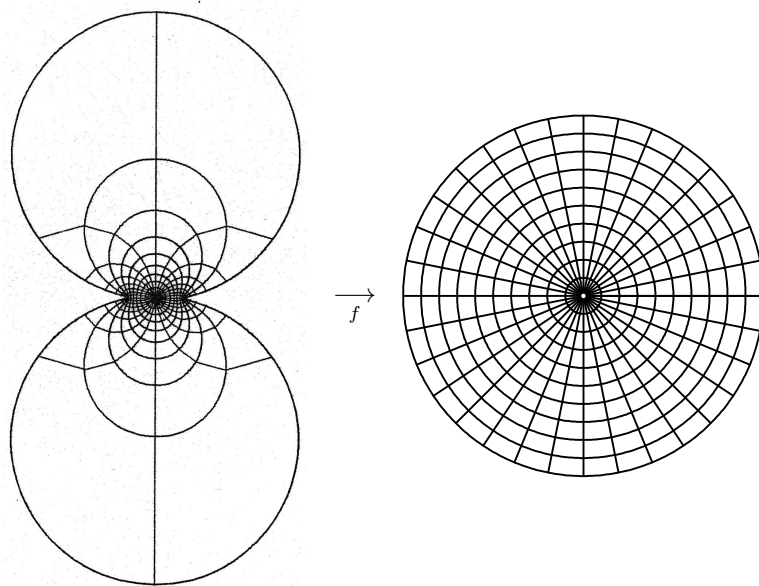


FIG. 2.3. Example for Riemann's mapping theorem for the domain Ω defined by $r(\phi) \leq (1 - 0.99 \cos^2 \phi)^{1/2}$; (computed by M. Gutknecht 18.12.1975)

on the boundary. (“Zwei gegebene einfach zusammenhängende ebene Flächen können stets so aufeinander bezogen werden, dass jedem Punkte der einen Ein mit ihm stetig fortrückender Punkt der anderen entspricht und ihre entsprechenden kleinsten Theile ähnlich sind;...”; Riemann [36], 1851, §21, *Werke* p. 40).

Riemann's proof of the mapping theorem: Riemann sketched a dubious proof of this audacious vision in the last two paragraphs of his thesis [36]. Some years later, in [37], he explained his ideas “in etwas veränderter Form”. The theorem is illustrated in Fig. 2.3, where Ω is an “inverse ellipse” defined by the radius $r(\phi) \leq (1 - 0.99 \cos^2 \phi)^{1/2}$; for z_0 we choose the origin. We then place a logarithm $\log(z - z_0) = \log r + i\phi$ at the point z_0 . The level curves of its real part $u_0(x, y)$ are the concentric circles around z_0 ; the level curves of the imaginary part, orthogonal to the first ones, are the star-shaped rays out of z_0 . The problem is that the boundary of Ω is normally not a level curve of u_0 . We postulate that there exists an everywhere harmonic function $u_1(x, y)$ such that $u_1(x, y) = u_0(x, y)$ on the boundary $\partial\Omega$. The function $u(x, y) := u_0(x, y) - u_1(x, y)$ will be harmonic in Ω with the exception of the point z_0 , where we have the logarithmic singularity, and it will be zero on $\partial\Omega$. By solving the differential equations (2.4), we complete $u(x, y)$ to a complex function $u(x, y) + iv(x, y)$. The exponential function of this will then map z_0 to the origin and the boundary of Ω to the boundary of the circle. Q.E.D.

These ideas of Riemann left much to do for scientists over the next century, clarifying the above proof, clarifying regularity conditions, more general domains, and better computational algorithms; see the last chapter of [23], Chapters 16 and 17 of Henrici's trilogy [19], and Gutknecht [18]. The main obstacle was the existence of the function u_1 , which we will discuss in the next subsection.

2.2. Dirichlet's Principle. Problem. Let Ω be a bounded domain and F an

arbitrary function defined on the boundary $\partial\Omega$. Find a function $w(x, y)$ with $\Delta w = 0$, such that $w = F$ for $(x, y) \in \partial\Omega$ (see Fig. 2.4 for an example).

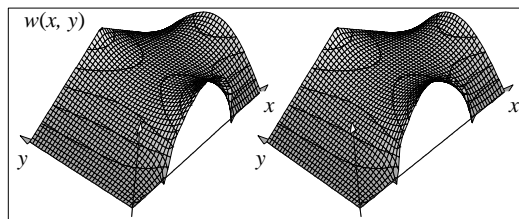


FIG. 2.4. Stereogram of the solution of $\Delta w = 0$, $0 \leq x \leq \Phi$, $\Phi = 1.618\dots$, $0 \leq y \leq 1$, $F(x, 0) = \sqrt{(-2 \cos \frac{2\pi x}{L})^+}$, $F(x, 1) = \min(\frac{\pi^2 x}{4L}, \frac{\pi^2(L-x)}{4L})$, $F(0, y) = 0$, $F(\Phi, y) = \sin \pi y$.

Riemann's solution [37]. This is based on the fact that the generalization of the variational problem (1.7) to higher dimensions is

$$(2.7) \quad \iint_{\Omega} \left(\frac{1}{2} \left(\left(\frac{\partial w}{\partial x} \right)^2 + \left(\frac{\partial w}{\partial y} \right)^2 \right) - f \cdot w \right) dx dy \longrightarrow \min \Rightarrow \boxed{-\frac{\partial^2 w}{\partial x^2} - \frac{\partial^2 w}{\partial y^2} = f.}$$

This relation had been discovered independently by Gauss [16], W. Thomson [29] and Dirichlet (in his lectures, which Riemann attended). The proof is precisely the same as above. Slightly more complicated is the step from (1.12) to (1.13): for one term we have partial integration with respect to x , which is straightforward, since the integration dx is inside the integration dy . For the second term we have partial integration with respect to y ; before doing this, we switch the order of integration:

“Hierzu kann in vielen Fällen ... ein Princip dienen, welches Dirichlet zur Lösung dieser Aufgabe für eine der Laplace'schen partiellen Differentialgleichung genügende Function ... in seinen Vorlesungen ... seit einer Reihe von Jahren zu geben pflegt.” (Riemann 1857, Werke p. 97)

Riemann concludes that for all functions defined on Ω with the prescribed boundary values F , the integral

$$(2.8) \quad J(w) = \iint_{\Omega} \frac{1}{2} \left(\left(\frac{\partial w}{\partial x} \right)^2 + \left(\frac{\partial w}{\partial y} \right)^2 \right) dx dy \quad \text{is always} > 0.$$

Choose among these functions the one for which this integral has the smallest possible value !

In other words, we pass through the Euler-Lagrange highway (1.11) \Rightarrow (1.14) **in the opposite direction** (1.11) \Leftarrow (1.14). In contrast to “Euler's world”, here the solution of the differential equation (1.14) is hopeless, whereas the variational problem (1.11) appears “trivial”. From here originates the name “Dirichlet's principle” and “Dirichlet boundary conditions”.

Weierstrass's Critique: Soon after, some people started to mistrust this “Brave new world”, the most serious critics came from Weierstrass [48] with the counter example

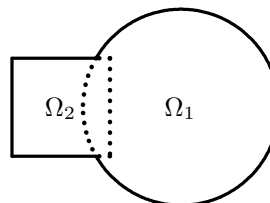
$$(2.9) \quad \int_{-1}^1 (x \cdot y')^2 dx \longrightarrow \min \left| \begin{array}{l} y = \frac{a+b}{2} + \frac{b-a}{2} \frac{\arctan \frac{x}{\epsilon}}{\arctan \frac{1}{\epsilon}} \\ y(-1) = a, \quad y(1) = b \end{array} \right| \quad \left| \begin{array}{c} \text{Graph of } y(x) \text{ showing a sharp transition from } a \text{ at } x=-1 \text{ to } b \text{ at } x=1. \end{array} \right.$$

The factor x in the integral allows y' to do anything close to the origin, and the solution of the problem becomes discontinuous (“Die Dirichlet’sche Schlussweise führt also in dem betrachteten Falle offenbar zu einem falschen Resultat.”)

F. Klein (*Entw. Math. 19. Jahrh.*, p. 264) reports that Riemann had answered to Weierstrass “my theorems remain nevertheless true” (meine Existenztheoreme sind trotzdem richtig), and that Helmholtz declared “for us physicists Dirichlet’s principle remains a proof” (Für uns Physiker bleibt das Dirichletsche Prinzip ein Beweis).

Proof without Dirichlet Principle: The majority of mathematicians, however, abandoned Riemann’s idea and started to look for alternative methods of proof. One of these was the invention of the *alternating method* by H.A. Schwarz (1870, Crelle 74, 1872). It is based on the fact that for rectangles, as well as circles, the existence of the solution is assured by Fourier series methods. A more complicated domain is represented as an overlapping union of such simple domains:

(2.10) Solve alternatively in Ω_1 and Ω_2 ;
 new boundary values on dotted curves;
 prove that iteration converges;
 add third, fourth domain etc.



Revival of Dirichlet’s Principle by Hilbert: Precisely half a century after Riemann’s thesis, D. Hilbert undertook the task of putting Dirichlet’s principle on a new basis [20, 21]. Certainly, Weierstrass’ counter-example is embarrassing, but the integral in (2.9) is not the same as in (2.8). So Hilbert managed, in an extraordinary *tour de force*, to establish a way of proving the existence of u directly from the properties of the integral (2.8). A couple of theses written under Hilbert’s supervision, including Courant’s [8] and [9], then improved this theory by making it simpler and more complete. These discussions became another motivation for W. Ritz for developing *his* method.

“Das Dirichletsche Prinzip verdankte seinen Ruhm der anziehenden Einfachheit seiner mathematischen Grundidee, dem unleugbaren Reichtum der möglichen Anwendungen ... und der ihm innewohnenden Überzeugungskraft.” (D. Hilbert)

“Mittlerweile war das verachtete und scheinote *Dirichletsche* Prinzip durch Hilbert wieder zum Leben erweckt worden;...” (Hurwitz-Courant [23], p. 392)

3. The Elastic Plate. The main motivation for Ritz was the announcement of the *Prix Vaillant* for 1907 of the Academy of Science in Paris (see Fig. 3.1), published in vol. 53 of *Journal für Mathematik und Physik*, p. 65. This announcement was sent by his friend Paul Ehrenfest to W. Ritz on a post card, in order for the “Scheusaltheorie” of Ritz’s thesis to find a new application. The deformation of an elastic plate under an external force was a very difficult problem at that time; it was first considered by Sophie Germaine in several articles (1811/13/15)⁶, to which Lagrange and Poisson had added corrections and improvements. The definitive breakthrough was achieved in a long article by Kirchhoff [24] (1850) in form of the differential equation

$$(3.1) \quad \Delta \Delta w = \frac{\partial^4 w}{\partial x^4} + 2 \frac{\partial^4 w}{\partial x^2 \partial y^2} + \frac{\partial^4 w}{\partial y^4} = f(x, y),$$

⁶This research was motivated by a visit of Chladni in Paris in 1808.

to which suitable boundary conditions must be added. If the plate is assumed to be clamped on all sides, we would have

$$(3.2) \quad w = 0 \quad \text{and} \quad \frac{\partial w}{\partial n} = 0 \quad \text{on } \partial\Omega.$$

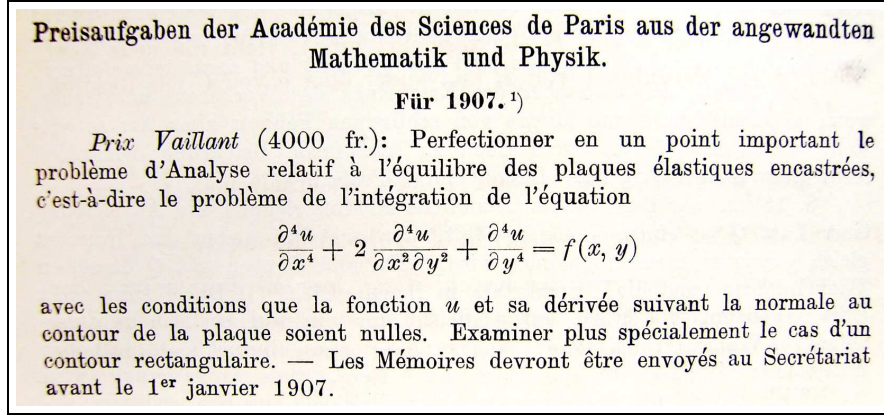


FIG. 3.1. Announcement for Prix Vaillant in *Journal für Mathematik und Physik*

Ritz had worked with many such problems in his thesis, where he tried to explain the Balmer series in spectroscopy (1902); it therefore appeared to him that he had good chances to succeed in this competition.

4. Ritz's Treatment of the Plate and Dirichlet's Principle. We now come to the publication in the Crelle Journal [38]



through which Ritz achieved immortality. As did Riemann, we proceed on the Euler-Lagrange highway backwards and arrive, starting from (3.1), at the variational problem

$$(4.1) \quad J = \iint_{\Omega} \left(\frac{1}{2} \left(\frac{\partial^2 w}{\partial x^2} \right)^2 + 2 \left(\frac{\partial^2 w}{\partial x^2} \right) \left(\frac{\partial^2 w}{\partial y^2} \right) + \left(\frac{\partial^2 w}{\partial y^2} \right)^2 \right) - f \cdot w \, dx \, dy \longrightarrow \min$$

or

$$(4.2) \quad J = \iint_{\Omega} \left(\frac{1}{2} (\Delta w)^2 - f \cdot w \right) dx \, dy \longrightarrow \min .$$

The standard procedure (“wie man ohne weiteres einsieht”) transforms term by term the formula (4.1) to (3.1); we just have to perform each time *two* integrations by parts, so that the minus sign in (2.7) disappears again.

4.1. Ritz's Method. The main idea is the following: we choose a sequence of functions

$$(4.3) \quad \psi_1(x, y), \psi_2(x, y), \psi_3(x, y), \psi_4(x, y), \dots$$

and try to approximate the solution of (4.1) as a linear combination⁷

$$(4.4) \quad w(x, y) = a_1\psi_1(x, y) + a_2\psi_2(x, y) + \dots + a_m\psi_m(x, y)$$

with the coefficients a_1, a_2, \dots, a_m to be determined. The quality of the method depends, of course, on a good choice of these functions. For example, if we assume Ω to be a square, one choice discussed by Ritz is

$$(4.5) \quad \begin{aligned} \psi_1(x, y) &= (1 - x^2)^2(1 - y^2)^2 \\ \psi_2(x, y) &= (1 - x^2)^2(1 - y^2)^2(x^2 + y^2) \\ \psi_3(x, y) &= (1 - x^2)^2(1 - y^2)^2(x^4 + y^4) \\ \psi_4(x, y) &= (1 - x^2)^2(1 - y^2)^2x^2y^2 \\ \psi_5(x, y) &= (1 - x^2)^2(1 - y^2)^2(x^6 + y^6) \\ \psi_6(x, y) &= (1 - x^2)^2(1 - y^2)^2(x^4y^2 + x^2y^4) \dots \end{aligned}$$

Each of these ψ_i contains the factor $(1 - x^2)^2(1 - y^2)^2$ in order to ensure the validity of the boundary conditions (3.2) for any choice of the a_i . We suppose $f(x, y) = 1$ and conclude by symmetry that odd powers of x or y will not appear. For the same reason, the terms containing x^2 and y^2 will have the same coefficient, so we simplify our calculations by collecting such equal terms into one function ψ . Another good property of our basis is that, for $m \rightarrow \infty$, it contains all (symmetric) polynomials; therefore they are able, by Weierstrass' approximation theorem, to approach any (symmetric) solution function w . This will be important for establishing convergence of the method.

If we insert the expression (4.4) into (4.2), we obtain a *finite-dimensional* expression in a_1, a_2, \dots, a_m . By chance, J in (4.2) contains only quadratic and linear terms of w . Hence, if we multiply out all the terms, we obtain a *finite-dimensional quadratic function*

$$(4.6) \quad J_m = \frac{1}{2} \sum_{i,j=1}^m k_{ij} a_i a_j - \sum_{i=1}^m b_i a_i$$

with

$$(4.7) \quad k_{ij} = \int_{-1}^1 \int_{-1}^1 \Delta\psi_i \cdot \Delta\psi_j \, dx \, dy, \quad b_i = \int_{-1}^1 \int_{-1}^1 \psi_i \cdot f \, dx \, dy.$$

Differentiating (4.6) with respect to a_ℓ , for $\ell = 1, 2, \dots, m$, we find that J_m is minimal if

$$(4.8) \quad \sum_{j=1}^m k_{\ell j} a_j = b_\ell.$$

This is a linear system which, in theory, is easily solvable. In practice, however, the system is very tedious to solve: for $m = 6$, obtaining just one of the 21 k -values would

⁷See the memorial plate in honor of Ritz at the Collège des Creusets, Sion, Switzerland.

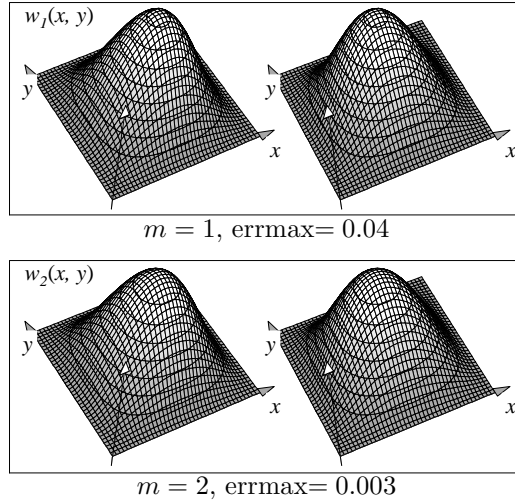


FIG. 4.1. Solutions of the plate problem for $m = 1$ and $m = 2$

require us to compute

$$\begin{aligned}
\Delta\psi_6 \cdot \Delta\psi_6 = & 8x^2(1-y^2)^2(x^4y^2 + x^2y^4) \\
& -8(1-x^2)(1-y^2)^2(4x^3y^2 + 2xy^4)x \\
& -4(1-x^2)(1-y^2)^2(x^4y^2 + x^2y^4) \\
& + (1-x^2)^2(1-y^2)^2(12x^2y^2 + 2y^4) \\
& +8(1-x^2)^2y^2(x^4y^2 + x^2y^4) \\
& -8(1-x^2)^2(1-y^2)(2x^4y + 4x^2y^3)y \\
& -4(1-x^2)^2(1-y^2)(x^4y^2 + x^2y^4) \\
& + (1-x^2)^2(1-y^2)^2(2x^4 + 12x^2y^2)
\end{aligned}$$

and then $\int_{-1}^1 \int_{-1}^1 \Delta\psi_6 \cdot \Delta\psi_6 dx dy = \frac{3052404736}{6898776885}$. We can understand why Ritz did not use this particular basis to demonstrate his method numerically. We, however, have modern computers and find the system

$$(4.9) \quad \begin{array}{cccccc|c}
53.49878 & 9.72705 & 2.13385 & 0.54039 & 0.63739 & 0.30484 & 1.13778 \\
9.72705 & 21.94822 & 10.91868 & 1.74588 & 5.99358 & 1.25965 & 0.32508 \\
2.13385 & 10.91868 & 9.63314 & 0.90947 & 7.10042 & 1.01770 & 0.10836 \\
0.54039 & 1.74588 & 0.90947 & 0.47237 & 0.51468 & 0.39553 & 0.02322 \\
0.63739 & 5.99358 & 7.10042 & 0.51468 & 6.32258 & 0.73117 & 0.04925 \\
0.30484 & 1.25965 & 1.01770 & 0.39553 & 0.73117 & 0.44246 & 0.01548
\end{array}$$

If we set $m = 1$, only $k_{11}a_1 = b_1$ has to be considered and we obtain by a simple division

$$(4.10) \quad a_1 = 1.13778/53.49878 = 0.021267 .$$

The solution $w = a_1\psi_1$ is drawn in the first picture of Fig. 4.1 and has an error of 4%. If we also include polynomial terms of degree 6, i.e., if we set $m = 2$, we have a linear system of two equations to solve and obtain

$$(4.11) \quad a_1 = 0.02020, \quad a_2 = 0.00586,$$

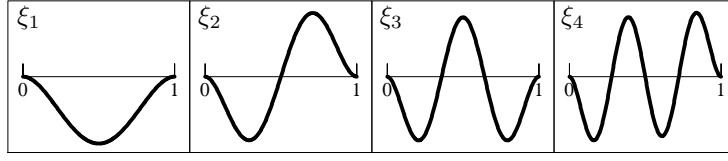


FIG. 4.2. One-dimensional elastic curves

which reduces the error to 0.3% (lower picture of Fig. 4.1). Taking into account also 8th degree terms ($m = 4$) leads to a relative error of 10^{-5} with precisely the same graphical representation as for $m = 2$. For $m = 6$ the solutions are

$$(4.12) \quad \begin{aligned} a_1 &= 0.02025, & a_2 &= 0.00521, & a_3 &= 0.00028, \\ a_4 &= 0.00612, & a_5 &= -0.00002, & a_6 &= 0.00012. \end{aligned}$$

4.2. A Basis Using One-Dimensional Elastic Curves. As satisfactory as the above numerical results are, they are not suitable for hand calculations. Instead of these polynomials, Ritz suggests using the normal modes of the clamped rod:

$$(4.13) \quad \xi^{(4)} = K^4 \cdot \xi \quad \xi(0) = \xi(1) = 0, \quad \xi'(0) = \xi'(1) = 0.$$

Ritz cites Lord Rayleigh [41], *Theory of Sound*, p.174 for the formulas; however, they had already been studied extensively by Euler, in an *Additamentum I, De curvis elasticis* of [12], §68–97.

Standard methods for differential equations with constant coefficients lead, together with the first three boundary conditions, to the formula

$$(4.14) \quad \xi_n(x) = \cos K_n x - \cosh K_n x - (\sin K_n x - \sinh K_n x) \frac{\cos K_n - \cosh K_n}{\sin K_n - \sinh K_n}.$$

The condition $\xi'(1) = 0$ then gives $\cos K_n \cosh K_n = 1$, which determines the discrete values of K_n as follows (see Fig. 4.2),

$$(4.15) \quad K_1 = 4.7300, \quad K_2 = 7.8532, \quad K_3 = 10.9956, \quad K_4 = 14.1372, \dots$$

Because $\cosh K$ tends to infinity very rapidly, K_n will also quickly approach the roots of the cosine, i.e. $K_n \approx (n + \frac{1}{2})\pi$.

Taking again into account the symmetry of the solution, we use the basis

$$\begin{aligned} \psi_1(x, y) &= \xi_1(x)\xi_1(y), & \psi_2(x, y) &= \xi_1(x)\xi_3(y) + \xi_3(x)\xi_1(y), \\ \psi_3(x, y) &= \xi_3(x)\xi_3(y), & \psi_4(x, y) &= \xi_1(x)\xi_5(y) + \xi_5(x)\xi_1(y), \dots \end{aligned}$$

The advantages of this basis are that

1. we only have one defining formula (4.14) for every n ,
2. using integration by parts, the condition (4.13) leads to easy formulas for the integrals (4.7), and
3. the linear system (4.8) becomes strongly diagonal dominant.

$l \cdot 431,5 = \underline{651,8} a_{11} - 239,2 a_{13} + 94,7 a_{33} - 188 a_{15} + 148 a_{35} + 58 a_{55}$
$l \cdot 188,5 = -119,6 a_{11} + \underline{8867} a_{13} - 961 a_{33} - 226 a_{15} - 515 a_{35} + 186 a_{55}$
$l \cdot 82,7 = 94,6 a_{11} - 1922 a_{13} + \underline{24390} a_{33} + 474 a_{15} - 4820 a_{35} + 592 a_{55}$
$l \cdot 120,0 = -93,8 a_{11} - 226 a_{13} + 237 a_{33} + \underline{48100} a_{15} - 2370 a_{35} - 2015 a_{55}$
$l \cdot 52,5 = 74,3 a_{11} - 515 a_{13} - 2410 a_{33} - 2370 a_{15} + \underline{78600} a_{35} - 6420 a_{55}$
$l \cdot 33,5 = 58 a_{11} + 372 a_{13} + 592 a_{33} - 4030 a_{15} - 12840 a_{35} + \underline{158500} a_{55}$

It can thus be solved easily (“der Rechenschieber angewandt werden kann...”) by a sort of “Gauss-Seidel” iteration. (“Eine direkte Lösung durch Determinanten würde 5stellige Logarithmentafeln erfordern”). In this way, Ritz obtained the solutions

$$\begin{aligned} w_1(x, y) &= 0.6620\xi_1(x)\xi_1(y), \\ w_2(x, y) &= 0.6727\xi_1(x)\xi_1(y) + 0.0307(\xi_1(x)\xi_3(y) \\ &\quad + \xi_3(x)\xi_1(y)) + 0.0031\xi_3(x)\xi_3(y), \\ w_3(x, y) &= 0.6740\xi_1(x)\xi_1(y) + 0.0380(\xi_1(x)\xi_3(y) + \xi_3(x)\xi_1(y)) \\ &\quad + 0.0032\xi_3(x)\xi_3(y) + 0.0040(\xi_1(x)\xi_5(y) + \xi_5(x)\xi_1(y)) \\ &\quad + 0.0004(\xi_3(x)\xi_5(y) + \xi_5(x)\xi_3(y)) + 0.0000\xi_5(x)\xi_5(y). \end{aligned}$$

whose graphical representation is not so different from those in Fig. 4.1.

Theorem of Existence and Convergence. Based on

- a careful study of speed of convergence based on the asymptotic values of K_n , allowing the exchange of differentiations and limits;
- Weierstrass’s approximation theorem;
- a modification of a lemma from Hilbert’s proof (1901),

Ritz managed, in paragraphs 3, 4, and 5 of his paper, to prove rigorously that the $w_m(x, y)$ converge to a function $w(x, y)$ that solves the minimization problem.

4.3. Dirichlet’s Principle. In the second part of [38], Ritz applies his proofs to Dirichlet’s problem (see Section 2.2)

$$\Delta w = 0 \quad w|_{\partial\Omega} = F .$$

In order to have the boundary values equal to 0, we subtract F from w and obtain a problem of the type

$$(4.16) \quad -\Delta w = f \quad w|_{\partial\Omega} = 0 .$$

Inserting (4.4) into (2.7), we obtain this time

$$(4.17) \quad k_{ij} = \int_{-1}^1 \int_{-1}^1 \left(\frac{\partial\psi_i}{\partial x} \frac{\partial\psi_j}{\partial x} + \frac{\partial\psi_i}{\partial y} \frac{\partial\psi_j}{\partial y} \right) dx dy , \quad b_i = \int_{-1}^1 \int_{-1}^1 \psi_i \cdot f dx dy .$$

instead of (4.7). Ritz then proceeds to prove once again the existence and convergence of his method, but he does not show numerical examples; for the square they would have been too simple.

Let us show them here, by modifying the basis functions according to the new boundary condition as

$$\begin{aligned} \psi_1(x, y) &= (1 - x^2)(1 - y^2) \\ \psi_2(x, y) &= (1 - x^2)(1 - y^2)(x^2 + y^2) \\ \psi_3(x, y) &= (1 - x^2)(1 - y^2)(x^4 + y^4) \\ \psi_4(x, y) &= (1 - x^2)(1 - y^2)x^2y^2 \\ \psi_5(x, y) &= (1 - x^2)(1 - y^2)(x^6 + y^6) \\ \psi_6(x, y) &= (1 - x^2)(1 - y^2)(x^4y^2 + x^2y^4) \dots \end{aligned}$$

and obtain, proceeding as above, the solutions

$$(4.18) \quad \begin{aligned} \text{for } m = 1 : \quad a_1 &= 0.3125 , \\ \text{for } m = 2 : \quad a_1 &= 0.2922 \quad a_2 = 0.05923 . \end{aligned}$$

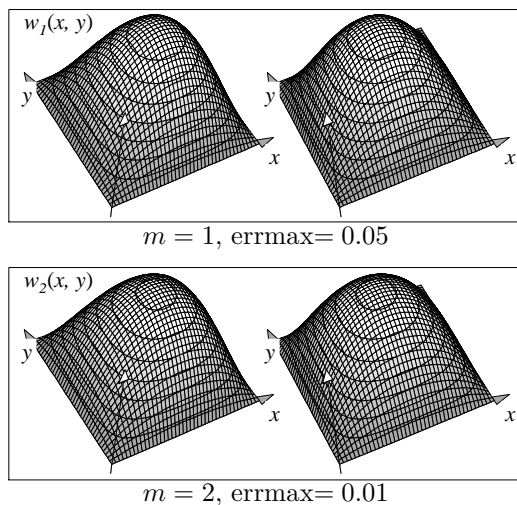


FIG. 4.3. Solutions for Dirichlet's principle

The solutions are drawn in Fig. 4.3. Again, the higher order terms have little effect on the graphic representation of the function. One observes that for each increase of the degree by 2, the error decreases by a factor 4.

5. Ritz Computes Chladni Figures. In 1787, Ernst Florence Friedrich Chladni, a musician and physicist from Wittenberg close to Leipzig, made an extraordinary discovery [5]: he noticed that when he tried to excite a metal plate with the bow of his violin, he could make sounds of different pitch, depending on where he touched the plate with the bow, see the experimental setup shown in the bottom left of Fig. 5.1. The plate itself was fixed only in the center, and when there was some dust or sand on the plate, for each pitch a beautiful pattern appeared. Chladni carefully collected all the figures he was able to create, and made drawings of each and every one, see Figure 5.1 on the right. These figures, now called Chladni figures after their inventor, attracted great attention among scientists and laymen alike, because of their intriguing beauty. We show in Fig. 5.2 more recent high tech experiments from Munich and San Diego, which recreate Chladni's experiments to very high accuracy. We invite the reader to compare those results with the historical drawings of Chladni to find similarities and differences.

5.1. The Mathematical Model for Chladni Figures. For the vibrating plate, the above equations (3.1) and (3.2) of Kirchhoff [24] (1850) are to be modified as follows. Chladni figures on a square plate correspond to eigenpairs (eigenvalues and corresponding eigenfunctions) of the bi-harmonic operator

$$(5.1) \quad \Delta^2 w = \lambda w \quad \text{in } \Omega := (-1, 1)^2,$$

with the free boundary conditions

$$(5.2) \quad \begin{aligned} \frac{\partial}{\partial x} \left(\frac{\partial^2 w}{\partial x^2} + (2 - \mu) \frac{\partial^2 w}{\partial y^2} \right) &= 0, & \frac{\partial^2 w}{\partial x^2} + \mu \frac{\partial^2 w}{\partial y^2} &= 0, & x &= \{-1, 1\}, \\ \frac{\partial}{\partial y} \left(\frac{\partial^2 w}{\partial y^2} + (2 - \mu) \frac{\partial^2 w}{\partial x^2} \right) &= 0, & \frac{\partial^2 w}{\partial y^2} + \mu \frac{\partial^2 w}{\partial x^2} &= 0, & y &= \{-1, 1\}, \end{aligned}$$

where μ here is the elasticity constant. Even before the correct mathematical model had been established, Wheatstone [49] had tried in 1833 to approximate Chladni

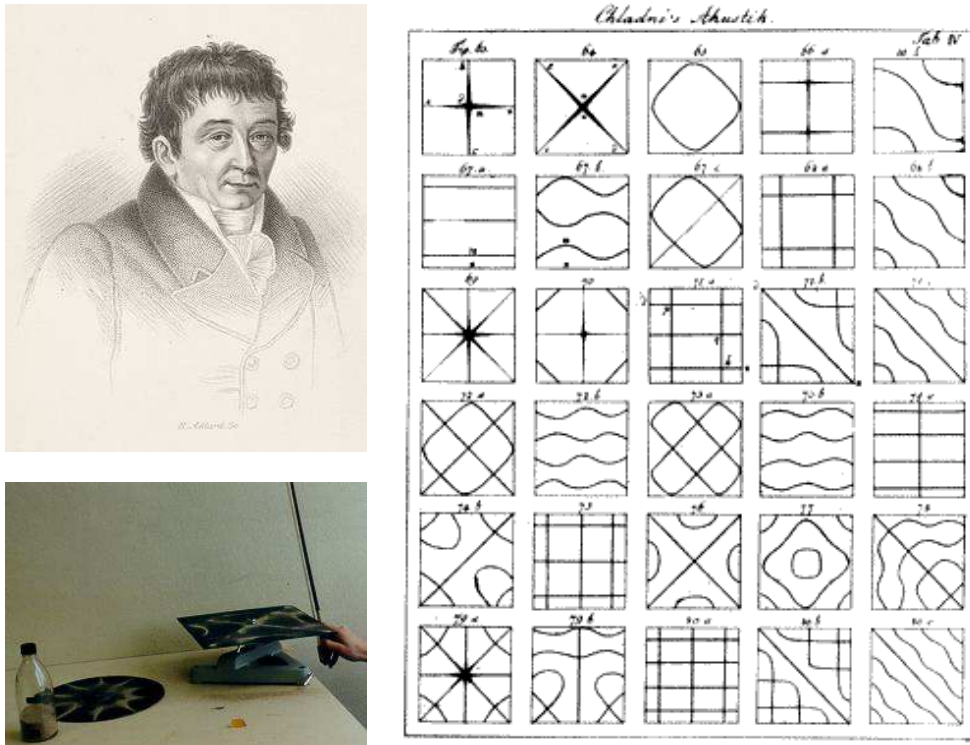


FIG. 5.1. Ernst Florens Friedrich Chladni and his famous experiment. On the right the drawings of Chladni of the figures he managed to create with his violin bow exciting an iron plate with some sand on it

figures using sine and cosine functions⁸. Kirchhoff, who came up with the correct mathematical model in 1850, also managed to solve the Chladni problem for the special case of a circular plate, which, due to symmetry, is much easier to handle. For other configurations, the partial differential eigenvalue problem (5.1) with the free boundary conditions (5.2) simply proved to be too difficult to solve⁹. There were even erroneous attempts: based on experiments performed by the master violin maker König [26] in 1864, which suggested that Chladni figures can only contain straight lines, Tanaka tried in 1887 to obtain solutions by integration starting from straight lines¹⁰. In the case of clamped boundaries, the problem greatly simplifies, and Voigt [47] found the general solution in 1893 for a rectangular plate with two or four clamped boundaries by elementary integration. Toward the end of the 19th century, the great expert in sound, John William Strutt, later Baron Rayleigh, summarized the situation in [41]: “The Problem of a rectangular plate, whose edges are free, is one of great

⁸Ritz: “... , dass es sich hier nur um einen in besonderen Fällen anwendbaren Kunstgriff handelt”

⁹Ritz: “Die von Kirchhoff erhaltene partielle Differentialgleichung ist vierter Ordnung, und es müssen am Rande zwei Differentialausdrücke dritter und zweiter Ordnung verschwinden, die von einer Elastizitätskonstante[n] abhängen. Die grosse hieraus sich ergebende Komplikation des Problems erklärt es hinreichend, dass die Lösung bis jetzt nur im Falle des Kreises (Kirchhoff) gefunden wurde, wobei sich ein sehr befriedigender Anschluss an die Erfahrung ergab. Die Klangfiguren bestehen hier nur aus konzentrischen Kreisen und aus Radien.”

¹⁰Ritz: “... Tanaka glaubt, allgemeinere und strengere Formeln zu erhalten. Dies ist aber schon deswegen nicht der Fall, weil übersehen ist, dass eine Randbedingung die Lösung gar nicht bestimmt.”

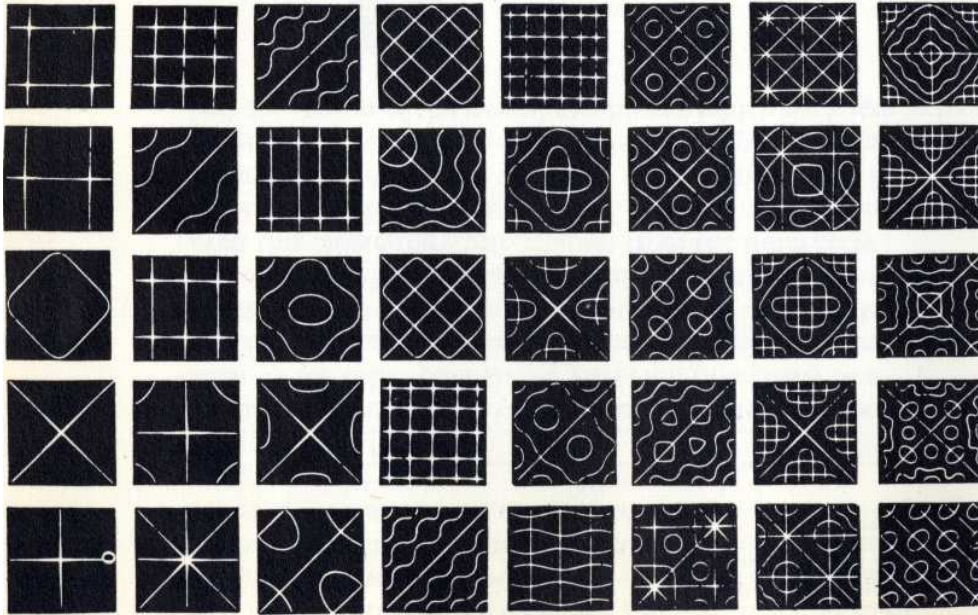


FIG. 5.2. Creation of Chladni figures in a modern laboratory environment in Munich and San Diego

difficulty, and has for the most part resisted attack”.

5.2. Ritz’s Computation of Chladni Figures. In his second groundbreaking paper [39], Walther Ritz presents and analyzes his own method in order to compute Chladni figures: instead of trying to solve the partial differential eigenvalue problem (5.1,5.2) directly, he proposes to use the principle of energy minimization, from which the equations were derived¹¹, and thus he considers the functional

$$(5.3) J(w) := \int_{-1}^1 \int_{-1}^1 \left[\left(\frac{\partial^2 w}{\partial x^2} \right)^2 + \left(\frac{\partial^2 w}{\partial y^2} \right)^2 + 2\mu \frac{\partial^2 w}{\partial x^2} \frac{\partial^2 w}{\partial y^2} + 2(1 - \mu) \left(\frac{\partial^2 w}{\partial x \partial y} \right)^2 \right].$$

¹¹Ritz: “Das wesentliche der neuen Methode besteht darin, dass nicht von den Differentialgleichungen und Randbedingungen des Problems, sondern *direkt vom Prinzip der kleinsten Wirkung* ausgegangen wird, aus welchem ja durch Variation jene Gleichungen und Bedingungen gewonnen werden können.”

According to the minimization principle, the solution w of (5.1,5.2) is a minimum of the constrained problem

$$(5.4) \quad J(w) \rightarrow \min, \quad \int_{-1}^1 \int_{-1}^1 w^2 dx dy = \text{const},$$

and from this minimization problem, one can obtain again the partial differential eigenvalue problem simply using again the central highway of variational calculus.

Even though Ritz explains his new method on the concrete example of Chladni figures on a square plate, he points out that his new method is completely general, and could be applied to plates of arbitrary shapes, provided the basis functions are well chosen¹². The fundamental idea of Ritz's new method is to search for an approximate solution of the problem as a combination of well chosen so-called coordinate functions ("Grundfunktionen"), of the form

$$(5.5) \quad w_s = \sum_{m=0}^s A_m w_m(x, y).$$

If a simple solution is sought, for example the lowest pitch, Ritz points out that one could simply choose polynomials for the basis functions $w_m(x, y)$ ¹³. For higher eigenmodes, a better choice that leads to more accuracy is to use the same coordinate functions as in the previous section¹⁴:

$$\begin{aligned} w_{mn} &= u_m(x)u_n(y) + u_m(y)u_n(x), \\ w'_{mn} &= u_m(x)u_n(y) - u_m(y)u_n(x), \end{aligned}$$

where $u_m(x)$ are the known eigenfunctions of a free one dimensional bar,

$$\frac{d^4 u_m}{dx^4} = k_m^4 u_m, \quad \text{with } \frac{d^2 u_m}{dx^2} = 0, \frac{d^3 u_m}{dx^3} = 0 \text{ at } x = \{-1, 1\}.$$

These conditions lead to the functions

$$(5.6) \quad u_m = \begin{cases} \frac{\cosh k_m \cos k_m x + \cos k_m \cosh k_m x}{\sqrt{\cosh^2 k_m + \cos^2 k_m}}, & \tan k_m + \tanh k_m = 0, \quad m \text{ even}, \\ \frac{\sinh k_m \sin k_m x + \sin k_m \sinh k_m x}{\sqrt{\sinh^2 k_m - \sin^2 k_m}}, & \tan k_m - \tanh k_m = 0, \quad m \text{ odd}. \end{cases}$$

Ritz approximation of the solution then takes the form

$$(5.7) \quad w_s := \sum_{m=0}^s \sum_{n=0}^s A_{mn} u_m(x) u_n(y).$$

¹²Ritz: "Im folgenden entwickle ich am Beispiel der quadratischen Platten mit freien Rändern eine neue Integrationsmethode, die ohne wesentliche Änderungen auch auf rechteckige Platten angewandt werden kann, sei es mit freien, sei es auch mit teilweise oder ganz eingespannten oder gestützten Rändern. Theoretisch ist die Lösung in ähnlicher Weise sogar für eine beliebige Gestalt der Platte möglich; eine genaue Berechnung einer grösseren Anzahl von Klangfiguren, wie sie im folgenden für den klassischen Fall der quadratischen Scheibe durchgeführt ist, wird aber nur bei geeigneter Wahl der Grundfunktionen, nach welchen entwickelt wird, praktisch ausführbar."

¹³Ritz: "Für den Grundton, wofern grosse Genauigkeit nicht gefordert wird, führt das Verfahren für die meisten Platten durch den Ansatz von Polynomen zum Ziel."

¹⁴Ritz: "Sämtliche Eigentöne der Platte lassen sich bis auf einige Prozent darstellen durch die Formeln..."

In order to determine the coefficients A_{mn} , we again insert this solution into the functional (5.3), and we require that the resulting functional

$$(5.8) \int_{-1}^1 \int_{-1}^1 \left[\left(\frac{\partial^2 w_s}{\partial x^2} \right)^2 + \left(\frac{\partial^2 w_s}{\partial y^2} \right)^2 + 2\mu \frac{\partial^2 w_s}{\partial x^2} \frac{\partial^2 w_s}{\partial y^2} + 2(1-\mu) \left(\frac{\partial^2 w_s}{\partial x \partial y} \right)^2 \right] dx dy,$$

be minimal¹⁵ under the constraint

$$(5.9) \quad U(w_s) := \int_{-1}^1 \int_{-1}^1 w_s^2 dx dy = \text{const} =: C.$$

In this approximate problem, we only need to determine a finite number of coefficients A_{mn} .

In order to express the functional $J(w_s)$ in terms of the coefficients A_{mn} , we have to evaluate several integral terms. The first one is

$$(5.10) \quad \begin{aligned} \int_{-1}^1 \int_{-1}^1 \left(\frac{\partial^2 w_s}{\partial x^2} \right)^2 dx dy &= \int_{-1}^1 \int_{-1}^1 \left(\frac{\partial^2 \sum_{m,n} A_{mn} u_m(x) u_n(y)}{\partial x^2} \right)^2 dx dy \\ &= \sum_{m,n} \sum_{p,q} A_{mn} A_{pq} \underbrace{\int_{-1}^1 \int_{-1}^1 \frac{\partial^2 u_m(x)}{\partial x^2} u_n(y) \frac{\partial^2 u_p(x)}{\partial x^2} u_q(y) dx dy}_{c_{mnpq}^1 :=} \end{aligned}$$

Since u_n is known, it suffices to evaluate the integrals numerically to obtain the coefficients c_{mnpq}^1 . Similarly, one can also evaluate all the other terms in (5.8), to obtain

$$\begin{aligned} \int_{-1}^1 \int_{-1}^1 \left(\frac{\partial^2 w_s}{\partial y^2} \right)^2 dx dy &= \sum_{m,n} \sum_{p,q} A_{mn} A_{pq} c_{mnpq}^2, \\ \int_{-1}^1 \int_{-1}^1 2\mu \frac{\partial^2 w_s}{\partial x^2} \frac{\partial^2 w_s}{\partial y^2} dx dy &= \sum_{m,n} \sum_{p,q} A_{mn} A_{pq} c_{mnpq}^3, \\ \int_{-1}^1 \int_{-1}^1 (1-\mu) \left(\frac{\partial^2 w_s}{\partial x \partial y} \right)^2 dx dy &= \sum_{m,n} \sum_{p,q} A_{mn} A_{pq} c_{mnpq}^4, \\ \int_{-1}^1 \int_{-1}^1 w_s^2 dx dy &= \sum_{m,n} A_{mn}^2 \quad \text{by orthogonality,} \end{aligned}$$

where all the coefficients c_{mnpq}^2 , c_{mnpq}^3 and c_{mnpq}^4 are defined by concrete integrals, like in (5.10). Using a Lagrange multiplier λ , we now need to minimize

$$J(w_s) - \lambda(U(w_s) - C) \longrightarrow \min,$$

which is equivalent to minimizing the finite dimensional problem

$$(5.11) \quad J_s(\mathbf{a}) := \mathbf{a}^T \tilde{K} \mathbf{a} - \lambda(\mathbf{a}^T \mathbf{a} - C) \longrightarrow \min$$

¹⁵Ritz: "Es liegt nahe, als Masstab des Gesamtfehlers die Abweichung der potentiellen Energie von ihrem exakten Wert beim wirklichen Vorgang zu wählen"

with respect to \mathbf{a} , where we defined the vector

$$\mathbf{a} := [A_{00}, A_{01}, A_{10}, \dots],$$

and the matrix

$$\tilde{K} := \begin{bmatrix} \alpha_{00}^{00} & \alpha_{01}^{00} & \alpha_{10}^{00} & \dots \\ \alpha_{00}^{01} & \alpha_{01}^{01} & \alpha_{10}^{01} & \dots \\ \alpha_{00}^{10} & \alpha_{01}^{10} & \alpha_{10}^{10} & \dots \\ \vdots & \vdots & \vdots & \ddots \end{bmatrix},$$

with $\alpha_{mn}^{pq} := c_{mnpq}^1 + c_{mnpq}^2 + c_{mnpq}^3 + c_{mnpq}^4$ using Ritz's original notation.

In order to minimize (5.11), we compute the gradient with respect to \mathbf{a} and set it to zero, to obtain with $K := \frac{1}{2}(\tilde{K} + \tilde{K}^T)$

$$(5.12) \quad K\mathbf{a} = \lambda\mathbf{a},$$

a discrete eigenvalue problem. For each eigenvalue λ^ℓ , we get an eigenvector $\mathbf{a}^\ell = [A_{00}^\ell, A_{01}^\ell, \dots]$, and the corresponding eigenfunction

$$w_s^\ell = \sum_{m=0}^s \sum_{n=0}^s A_{mn}^\ell u_m(x) u_n(y).$$

Note the similarity with the underlying continuous eigenvalue problem

$$\Delta^2 w = \lambda w$$

we started with in (5.1).

The infinite partial differential eigenvalue problem has been reduced to the numerical evaluation of several integrals in order to obtain the matrix, and to the solution of an eigenvalue problem using a numerical linear algebra package; both tasks are easy nowadays, if one has a computer available. However, at the time of Ritz, no such tools were available, and Ritz had the daunting tasks of computing approximations of all the integrals by hand, and then solving the eigenvalue problem, also by hand. For the integrals, Ritz had to first approximate the one-dimensional basis functions (5.6), which contain in their definition the solution of transcendental equations. Here he notices that only the first few values of k_m need to be computed¹⁶, since they quickly become close to $m\pi/2 - \pi/4$, and he then approximates the basis functions for even m by¹⁷

$$u_m = \cos\left(\frac{m}{2} - \frac{1}{4}\right)\pi x + \frac{(-1)^{\frac{m}{2}} \cosh\left(\frac{m}{2} - \frac{1}{4}\right)\pi x}{\sqrt{2} \cosh\left(\frac{m}{2} - \frac{1}{4}\right)\pi},$$

(by neglecting $\cos k_m$ in the denominator of the defining formula (5.6) compared to $\cosh k_m$) and by a similar expression for odd m . Further approximations are needed throughout (“Begnügt man sich mit vier genauen Ziffern . . .”, “Mit einer Genauigkeit von mindestens 2 Prozent . . .”, “Zur Vereinfachung wird man die aus der Symmetrie der Lösung sich ergebenden Beziehungen zwischen den A_{mn} sogleich einführen. . .”) before Ritz finally transforms the theoretical eigenvalue problem

¹⁶Because $\tanh k_m$ approaches very rapidly 1; Ritz: “Die Wurzeln von $\tan k_m + \tanh k_m = 0$ unterscheiden sich nur wenig von $m\pi/2 - \pi/4$.”

¹⁷Ritz: “Für $m > 2$ ist auf vier Stellen genau für gerade m ”

$$(53) \quad \begin{cases} 0 = (\alpha_{00}^{(00)} - \lambda_s) A_{00} + \alpha_{01}^{(00)} A_{01} + \alpha_{10}^{(00)} A_{10} + \dots + \alpha_{ss}^{(00)} A_{ss}, \\ 0 = \alpha_{00}^{(01)} A_{00} + (\alpha_{01}^{(01)} - \lambda_s) A_{01} + \alpha_{10}^{(01)} A_{10} + \dots + \alpha_{ss}^{(01)} A_{ss}, \\ \dots \\ 0 = \alpha_{00}^{(ss)} A_{00} + \alpha_{01}^{(ss)} A_{01} + \alpha_{10}^{(ss)} A_{10} + \dots + (\alpha_{ss}^{(ss)} - \lambda_s) A_{ss}. \end{cases}$$

into the concrete numerical one, which contains six coefficients A_0, \dots, A_5 to be determined for a particular symmetry in the eigenfunction:

$u_m(y) = v_m$ gesetzt und

$$w = A_0 u_1 v_1 + A_1 (u_1 v_3 + v_1 u_3) + A_2 u_3 v_3 + A_3 (u_1 v_5 + u_5 v_1) \\ + A_4 (u_3 v_5 + u_5 v_3) + A_5 u_5 v_5.$$

Das System (53) wird hier

$$(54) \quad \begin{cases} 0 = (13,95 - \lambda) A_0 - 32,08 A_1 + 18,60 A_2 + 32,08 A_3 - 37,20 A_4 + 18,60 A_5, \\ 0 = -16,04 A_0 + (411,8 - \lambda) A_1 - 120,0 A_2 - 133,6 A_3 + 166,8 A_4 + 140 A_5, \\ 0 = +18,60 A_0 - 240,0 A_1 + (1686 - \lambda) A_2 - 218,0 A_3 - 1134 A_4 + 330 A_5, \\ 0 = +16,04 A_0 - 133,6 A_1 + 109,0 A_2 + (2945 - \lambda) A_3 - 424 A_4 + 179 A_5, \\ 0 = -18,6 A_0 + 166,8 A_1 - 567 A_2 - 424 A_3 + (6303 - \lambda) A_4 - 1437 A_5, \\ 0 = +18,6 A_0 + 280 A_1 - 330 A_2 + 358 A_3 - 2874 A_4 + (13674 - \lambda) A_5. \end{cases}$$

Nowadays, the symbolic calculator Maple can easily perform all these operations! We first define the solution of the transcendental functions defining k_m in (5.6) using the command

```
k:=m->if type(m,even) then
  fsolve(tan(x)+tanh(x)=0,x=m*Pi/2-Pi/4)
else
  fsolve(tan(x)-tanh(x)=0,x=(m-1/2)*Pi/2)
end if;
```

Having k_m available, we can now define the one dimensional eigenfunctions u_m used by Ritz to construct his coordinate functions, using the Maple command

```
u:=(m,x)->if m=0 then
  1/sqrt(2)
elif m=1 then
  sqrt(3/2)*x
elif type(m,even) then
  (cosh(k(m))*cos(k(m)*x)+cos(k(m))*cosh(k(m)*x))/
  sqrt((cosh(k(m)))^2+(cos(k(m)))^2)
else
  (sinh(k(m))*sin(k(m)*x)+sin(k(m))*sinh(k(m)*x))/
  sqrt((sinh(k(m)))^2-(sin(k(m)))^2)
end if;
```

We now simply have to evaluate numerically the integrals in order to obtain the matrix entries,

```
mu:=0.225;
for m from 1 to 5 by 2 do
  for n from 1 to 5 by 2 do
    for p from 1 to 5 by 2 do
      for q from 1 to 5 by 2 do
        K(m,n,p,q):=evalf(Int(Int(diff(u(m,x),x,x)*u(n,y)*diff(u(p,x),x,x)*u(q,y),
          x=-1..1),y=-1..1)
        +Int(Int(diff(u(n,y),y,y)*u(m,x)*diff(u(q,y),y,y)*u(p,x),
```

```

x=-1..1),y=-1..1)
+2*mu*Int(Int(diff(u(m,x),x,x)*u(n,y)*diff(u(q,y),y,y)*u(p,x),
x=-1..1),y=-1..1)
+2*(1-mu)*Int(Int(diff(u(m,x),x)*diff(u(n,y),y)*diff(u(p,x),x)*diff(u(q,y),y),
x=-1..1),y=-1..1));
od;
od;
od;
od;

```

and now we put the entries into a matrix¹⁸,

```

Ke:=matrix([[K(1,1,1,1),K(1,1,1,3)+K(1,1,3,1),K(1,1,3,3),K(1,1,1,5)
+K(1,1,5,1),K(1,1,3,5)+K(1,1,5,3),K(1,1,5,5)],
[(K(1,3,1,1)+K(3,1,1,1))/2,(K(1,3,1,3)+K(1,3,3,1)+K(3,1,3,1)+K(3,1,1,3))/2,
(K(1,3,3,3)+K(3,1,3,3))/2,(K(1,3,1,5)+K(1,3,5,1)+K(3,1,1,5)+K(3,1,5,1))/2,
(K(1,3,3,5)+K(1,3,5,3)+K(3,1,3,5)+K(3,1,5,3))/2,(K(1,3,5,5)+K(3,1,5,5))/2],
[K(3,3,1,1),K(3,3,1,3)+K(3,3,3,1),K(3,3,3,3),K(3,3,1,5)+K(3,3,5,1),
K(3,3,3,5)+K(3,3,5,3),K(3,3,5,5)],
[(K(1,5,1,1)+K(5,1,1,1))/2,(K(1,5,1,3)+K(1,5,3,1)+K(5,1,3,1)+K(5,1,1,3))/2,
(K(1,5,3,3)+K(5,1,3,3))/2,(K(1,5,1,5)+K(1,5,5,1)+K(5,1,1,5)+K(5,1,5,1))/2,
(K(1,5,3,5)+K(1,5,5,3)+K(5,1,3,5)+K(5,1,5,3))/2,(K(1,5,5,5)+K(5,1,5,5))/2],
[(K(5,3,1,1)+K(3,5,1,1))/2,(K(5,3,1,3)+K(5,3,3,1)+K(3,5,3,1)+K(3,5,1,3))/2,
(K(5,3,3,3)+K(3,5,3,3))/2,(K(5,3,1,5)+K(5,3,5,1)+K(3,5,1,5)+K(3,5,5,1))/2,
(K(5,3,3,5)+K(5,3,5,3)+K(3,5,3,5)+K(3,5,5,3))/2,(K(5,3,5,5)+K(3,5,5,5))/2],
[K(5,5,1,1),K(5,5,1,3)+K(5,5,3,1),K(5,5,3,3),K(5,5,1,5)+K(5,5,5,1),
K(5,5,3,5)+K(5,5,5,3),K(5,5,5,5)]];

```

This leads to the matrix

$$\begin{bmatrix} 13.95000000 & -32.21614500 & 18.60000000 & 32.21614504 & -37.20000000 & 18.60000002 \\ -16.10807250 & 406.1196695 & -119.9256010 & -133.5499734 & 172.0806028 & -83.55402336 \\ 18.60000000 & -239.8512019 & 1684.464298 & 217.9828187 & -1136.623991 & 329.5350388 \\ 16.10807252 & -133.5499735 & 108.9914093 & 2945.466708 & -427.1327599 & 179.2413502 \\ -18.60000000 & 172.0806029 & -568.3119956 & -427.1327599 & 6325.441344 & -1441.531357 \\ 18.60000002 & -167.1080467 & 329.5350388 & 358.4827005 & -2883.062714 & 13672.20612 \end{bmatrix}.$$

Comparing this result, computed accurately to about ten significant figures in Maple, with the results obtained by Ritz, who had to both approximate the roots of the transcendental functions, the functions themselves, and also intermediate results during the computation, we are stunned by the accuracy of Ritz's results: we show in Figure 5.3 the original results of Ritz, in red the digits which would need to be modified if Ritz had performed his computations to 10 digits accuracy, but with his approximations used for the functions, and in green the digits which would further need to be modified if no approximations to the functions had been made. We can see two likely misprints, the sign of the coefficient A_3 in the third equation, and the sign of the coefficient of A_2 in the last equation, and maybe one real error, the coefficient in front of A_5 in the second equation and A_1 in the last equation, where one should be the double of the other, and both are consistently incorrect by a large margin.

Once Ritz had obtained the discrete eigenvalue problem, he still needed to solve it in order to obtain eigenvalues and eigenvectors, which then allowed him to reconstruct the vibrational modes of the plate using the coordinate functions. At the time of Ritz, there were no numerical methods available to compute eigenvalues and eigenvectors of a matrix. And again Ritz proved to be far ahead of his time: instead of

¹⁸The particular combinations are due to the symmetries Ritz used.

$$\begin{aligned}
0 &= (13,95 - \lambda) A_0 - 32,08 A_1 + 18,60 A_2 + 32,08 A_3 - 37,20 A_4 + 18,60 A_5 \\
0 &= -16,04 A_0 + (411,8 - \lambda) A_1 - 120,0 A_2 - 133,6 A_3 + 166,8 A_4 + 140 A_5 \\
0 &= +18,60 A_0 - 240,0 A_1 + (1686 - \lambda) A_2 - 218,0 A_3 - 1134 A_4 + 330 A_5, \\
0 &= +16,04 A_0 - 133,6 A_1 + 109,0 A_2 + (2945 - \lambda) A_3 - 424 A_4 + 179 A_5, \\
0 &= -18,6 A_0 + 166,8 A_1 - 567 A_2 - 424 A_3 + (6303 - \lambda) A_4 - 1437 A_5, \\
0 &= +18,6 A_0 + 280 A_1 - 330 A_2 + 358 A_3 - 2874 A_4 + (13674 - \lambda) A_5.
\end{aligned}$$

FIG. 5.3. High Accuracy of the Results of Ritz

trying to compute the eigenvalues using the characteristic polynomial, a disastrous approach as we know today for numerical purposes, he invents an iterative method, an approach which has now become standard for eigenvalue problems and also linear systems, and for which very powerful methods are available today. Ritz first fixes one component of the eigenvector, namely $A_0 = 1$, and takes as an approximation to the eigenvalue the first diagonal entry. Then the remaining last 5 equations can be used to determine an approximation of the other components A_j of the eigenvector (“...setzen wir $A_0 = 1$, und in erster Annäherung $\lambda_0 = 13.95$. Dann ergeben die fünf letzten Gleichungen die übrigen A_i .”). But again, solving this linear system of five equations is too much work by hand, so Ritz proposes to just invert the diagonal of the matrix, what we would today call a Jacobi step (“Wir berechnen für die A_i eine erste Approximation, indem wir alle Glieder rechts vernachlässigen neben den Diagonalgliedern ...”). Having this approximation for the eigenvector, one can now compute a correction to the eigenvalue using the first equation, and according to Ritz, one or two successive iterations suffice in order to obtain about four digits of accuracy (“Ein oder zwei sukzessive Korrekturen genügen meist, um die vierte Stelle bis auf wenige Einheiten festzustellen.”). It is worthwhile to write this algorithm in today’s notation: solving the eigenvalue problem

$$K\mathbf{a} = \lambda\mathbf{a}$$

for $\mathbf{a} := (a_0, a_1, \dots, a_n)$ and λ is equivalent to solving the nonlinear system of equations

$$\mathbf{f}(\lambda, a_1, \dots, a_n) := K\mathbf{a} - \lambda\mathbf{a} = 0,$$

where we fixed, like Ritz, one component of the eigenvector, $a_0 := 1$. Ritz’s algorithm now starts with $\lambda^0 = 13.95$ and $(a_1^0, \dots, a_n^0) = \mathbf{0}$, and computes for iteration index $k = 0, 1, \dots$

$$(5.13) \quad f_j(\lambda^k, a_1^k, \dots, a_{j-1}^k, a_j^{k+1}, a_{j+1}^k, \dots, a_n^k) = 0, \quad j = 1, 2, \dots, n,$$

and then solves for a new approximation of the eigenvalue λ^{k+1}

$$f_0(\lambda^{k+1}, a_1^{k+1}, \dots, a_n^{k+1}) = 0.$$

Note that in each step of the algorithm, only a scalar linear equation needs to be solved. Implementing this method in Matlab gives the following short program:


```

K=[13.95  -32.08   18.60  32.08  -37.20  18.60
   -16.04  411.8   -120.0  -133.6  166.8  140
    18.60 -240.0  1686    -218.0 -1134   330
    16.04 -133.6   109.0  2945   -424   179
   -18.6   166.8   -567   -424  6303   -1437
    18.6   280    -330    358  -2874  13674];

Kr=K(2:end,2:end);           % last 5 equations
Kd=diag(Kr);                 % Jacobi splitting for Ritz'
Ko=Kr-diag(diag(Kr));       % eigenvalue iteration

la=K(1,1)                    % first eigenvalue and
a=zeros(1,size(K,2)); a(1)=1; % eigenvector approximation

for j=1:6                    % Ritz iteration
    b=-K(2:end,1);          % rhs of the system
    bj=b-Ko*a(2:end)';     % rhs for Jacobi step
    a(2:6)=bj./(Kd-la)     % Jacobi step
    la=K(1,1)+K(1,2:end)*a(2:6) % new eigenvalue approximation
end

```

whose output is

```

la =
    13.9500
a =
    1.0000    0.0403   -0.0111   -0.0055    0.0030   -0.0014
la =
    12.1388
a =
    1.0000    0.0342   -0.0038   -0.0027    0.0002   -0.0017
la =
    12.6565
a =
    1.0000    0.0387   -0.0061   -0.0036    0.0011   -0.0020
la =
    12.3996
a =
    1.0000    0.0374   -0.0049   -0.0032    0.0007   -0.0020
la =
    12.4972
a =
    1.0000    0.0380   -0.0053   -0.0034    0.0009   -0.0020
la =
    12.4525
a =
    1.0000    0.0378   -0.0051   -0.0033    0.0008   -0.0020
la =
    12.4711

```

we see that the dominant eigenvector component is the first one, corresponding to the diagonal element, and the eigenvalue indeed seems to converge. In order to obtain the exact eigenvalue of Ritz's matrix, we can use the Matlab commands

```

>> ev=sort(eig(K));
>> ev(1)
ans =
    12.4653

```

which first computes all eigenvalues, sorts them and then shows the first one. We see that about four iterations are needed in this case in order to determine the first three digits. It is also interesting to check if the approximations chosen by Ritz have a significant influence on the eigenvalues. We therefore computed the eigenvalues of Ritz's matrix, the matrix obtained when computing exactly using Ritz's approximations of the quantities in the matrix, and the exact matrix. We obtain for the eigenvalues

12.47	12.49	12.49
379.85	379.14	379.34
1579.79	1556.84	1559.28
2887.06	2899.82	2899.93
5969.67	5957.80	5961.32
14204.92	14233.73	14235.30

which shows that even with the small error in Ritz's computation in the elements, the approximation has little influence on the final result.

Having obtained these numerical results, all by hand calculations, Ritz now goes on to compare his results to the physical experiments performed by Chladni. We show in Fig. 5.4 the results for the first type of symmetry Ritz computed. We see that in case II, there is excellent agreement between the theoretical and experimental results (in case one trivially also). In case III, there were no measurements available. Case VI was not obtained by Chladni, and we see in Fig. 5.4 that fairly complicated patterns arise for the later cases. Ritz also used his method to estimate the pitch of the sound, and compared the result with the pitch measured by Chladni. Due to the enormous demand in computations however, Ritz only used his full method of approximation with several terms in the expansion for the first 9 notes, marked with a star in the original Table 5.5. For the remaining notes, he only uses the first term in the expansion. It is remarkable how well Ritz's results agree with the physical experiments of Chladni!

5.3. First Reactions to Ritz's Work in Western Europe. In Göttingen itself, where Ritz had spent the last years of his life and which was then *the* principal center of mathematical research, the importance of Ritz's invention was not immediately accepted. None of the many doctoral and habilitation theses written under Hilbert in these years on the Dirichlet Principle gives any reference to Ritz's work. The only exception is the paper by König [25], who just says that Ritz had simplified somewhere some lines of Hilbert's proof. Ten years later, we read in a footnote in the famous book of Hurwitz and Courant [23] (see Fig. 5.6) the following:

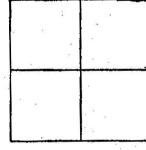
The actual, for this proof of existence *unimportant* [italics by us], construction of minimal sequences causes no problems in principle. For example if G is a finite domain, bounded by curves C without multiple points, we imagine that G is covered by a triangular grid T_j , whose triangles become smaller when j increases. We consider now only functions φ or $\phi = \varphi - S$, where the difference $\varphi - \frac{x}{x^2+y^2}$ in each Triangle of T_j is a linear function. The function ϕ_j is then the function associated with T_j , such that $D[\phi]$ has as small a value as possible. The requirement that $D[\phi] = \text{Min}$ is now a problem of minimization with a finite number of variables, namely the integrals, understood in dependence of the values of φ in the nodes of the triangular mesh; this problem can certainly be solved, as one can easily see, with linear equations. The fact that the functions so created

A. Lösungen, die in x und y ungerade und symmetrisch sind:

I. Grundton. $\lambda = 12,43 - 18,0\delta\mu$.

$$\begin{aligned} w = & \mathbf{u}_1 \mathbf{v}_1 + 0,0394(u_1 v_3 + v_1 u_3) \\ & - 0,0040 u_3 v_3 - 0,0034(u_1 v_5 + u_5 v_1) \\ & + 0,0011(u_3 v_5 + u_5 v_3) - 0,0019 u_5 v_5. \end{aligned}$$

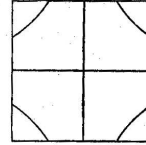
Fig. 1.



II. $\lambda = 378 - 57\delta\mu$.

$$\begin{aligned} w = & -0,075 u_1 v_1 + (\mathbf{u}_1 \mathbf{v}_3 + \mathbf{u}_3 \mathbf{v}_1) \\ & + 0,173 u_3 v_3 + 0,045(u_1 v_5 + u_5 v_1) \\ & - 0,015(u_3 v_5 + u_5 v_3) - 0,029 u_5 v_5. \end{aligned}$$

Fig. 2.



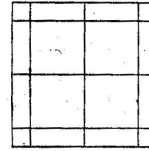
Es ist

y beob. :	0,530	0,578	0,630	0,690	0,752	0,819	0,893
x beob. :	0,9375	0,8750	0,8125	0,7500	0,6875	0,6250	0,5625
x ber. - x beob. :	-0,003	-0,002	0,000	-0,001	-0,0005	0,000	0,000

III. $\lambda = 1554$.

$$\begin{aligned} w = & 0,009 u_1 v_1 - 0,075(u_1 v_3 + v_1 u_3) \\ & + \mathbf{u}_3 \mathbf{v}_3 - 0,057(u_1 v_5 + u_5 v_1) \\ & + 0,121(u_3 v_5 + u_5 v_3) - 0,007 u_5 v_5. \end{aligned}$$

Fig. 3.



Messungen fehlen.

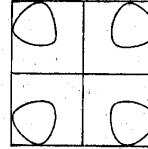
R.

20

IV. $\lambda = 2945$.

$$w = u_1 v_5 + u_5 v_1.$$

Fig. 4.



Diese Figur fehlt bei Chladni.

Fig. 5.

V. $\lambda = 6303$.

$$w = u_3 v_5 + u_5 v_3.$$

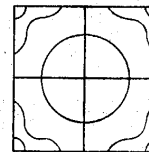


Fig. 6.

VI. $\lambda = 13674$.

$$w = u_5 v_5.$$

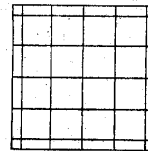


FIG. 5.4. Comparison done by Ritz, of his computational results with physical experiments, for the first few cases

Tabelle der Tönhöhen ($\mu = 0,225$).

HAUPTGLIEDER	λ	BER.	BOB.	HAUPTGLIEDER	λ	BER.	BOB.
$u_1 v_1$	12,43	G*	G	$u_3 v_4 \pm u_4 v_3$	3240	$g_3 +$	fis_3
$u_0 v_2 - v_0 u_2$	26,40	d^*	d	$u_5 v_2 \pm u_2 v_5$	3927	$a_3 +$	$gis_3 +$
$u_0 v_2 + v_0 u_2$	35,73	e^*	e	$u_4 v_4$	5480	$ais_3 +$	ais_3
$u_1 v_2 \pm u_2 v_1$	80,8	h^*	h	$u_0 v_6 - u_6 v_0$	5500	$c_4 -$	$-(^2)$
$u_0 v_3 \pm u_3 v_0$	237,1	$gis_1^* +$	$gis_1 +$	$u_3 v_5 - u_5 v_3$	5570	$c_4 -$	$ais_3 -$
$u_1 v_1$	266,0	$ais_1^* -$	$ais_1 -$	$u_0 v_6 + u_6 v_0$	5640	$c_4 -$	$-(^2)$
$u_1 v_3 - u_3 v_1$	316,1	h_1^*	h_1	$u_1 v_6 \pm v_1 u_6$	6036	$c_4 +$	$c_4 -$
$u_1 v_3 + u_3 v_1$	378	cis_2^*	cis_2	$u_3 v_5 + u_5 v_3$	6303	cis_4	$c_4 -$
$u_2 v_3 \pm u_3 v_2$	746	$fis_2^* +$	fis_2	$u_2 v_6 - u_6 v_2$	7310	$d_4 +$	$cis_4 +$
$u_0 v_4 - v_0 u_4$	886	gis_2	gis_2	$u_2 v_6 + u_6 v_2$	7840	$dis_4 -$	$d_4 -$
$u_0 v_4 + v_0 u_4$	941	$gis_2 +$	$gis_2 +$	$u_5 v_4 \pm u_4 v_5$	9030	e_4	dis_4
$u_1 v_4 \pm u_4 v_1$	1131	ais_2	$ais_2 -$	$u_6 v_3 \pm u_3 v_6$	10380	f_4	e_4
$u_3 v_3$	1554	$c_3 +$	c_3	$u_5 v_5$	13670	$g_4 +$	$fis_4 +$
$u_2 v_4 - u_4 v_2$	1702	$d_2 -$	cis_3	$u_6 v_4 - u_4 v_6$	13840	$g_4 +$	$g_4 +$
$u_2 v_4 + u_4 v_2$	2020	dis_3	d_3	$u_6 v_4 + u_4 v_6$	15120	$gis_4 +$	$g_4 +$
$u_0 v_5 \pm v_0 u_5$	2500	$f_3 -$	$f_3 -$	$u_6 v_5 \pm u_5 v_6$	20400	h_4	$ais_4 -$
$u_1 v_5 - v_1 u_5$	2713	fis_3	$fis_3 -$	$u_6 v_6$	28740	d_5	$-(^2)$
$u_1 v_5 + v_1 u_5$	2945	$fis_3 +$	$fis_3 (^1)$				

FIG. 5.5. Table of frequencies obtained by Ritz for the pitch of the sound of the plate, compared to the results measured by Chladni

¹⁾ Die wirkliche, für den bloßen Existenzbeweis unerhebliche Konstruktion solcher Minimalfolgen macht keinerlei prinzipielle Schwierigkeiten. Ist z. B. G ein ganz im Endlichen gelegener Bereich, begrenzt von Kurven C ohne mehrfache Punkte, so denken wir uns denselben mit einem noch von dem Index j abhängigen Dreiecksnetz T_j überdeckt, dessen Maschen mit wachsendem j immer enger werden. Wir betrachten nun nur solche Funktionen φ bzw. $\Phi = \varphi - S$, wo die Differenz $\varphi - \frac{x}{x^2 + y^2}$ in jedem Dreieck von T_j eine lineare Funktion wird. Unter Φ_j verstehen wir diejenige unter den so entstehenden zu T_j konstruierten Funktionen, für welche $D[\Phi]$ einen möglichst kleinen Wert erhält. Diese Forderung $D[\Phi] = \text{Min.}$ ist jetzt ein Problem eines Minimums einer Funktion von einer endlichen Anzahl von Variablen, nämlich des Integrals, aufgefaßt in seiner Abhängigkeit von den Werten von φ in den Eckpunkten der Dreieckseinteilung; dieses Problem ist gewiß lösbar, und zwar, wie leicht ersichtlich, mittels linearer Gleichungen. Daß die so entstehenden Funktionen Φ_j wirklich eine Minimalfolge bilden, folgt sofort aus der unschwer beweisbaren Tatsache, daß man jede zulässige Funktion Φ und deren Dirichletsches Integral mit Hilfe unserer Konstruktion bei hinreichend großem j beliebig genau approximieren kann.

FIG. 5.6. Footnote in the first edition of the book by Hurwitz and Courant [23]

швейцарского ученого Вальтера Ритца

FIG. 6.1. 'Shveitsarskogo uchenogo Val'tera Rittsa', as Timoshenko refers to the inventor of the method in his original publication in Russian

form a minimization sequence follows immediately from the easy to prove fact that one can approximate each admissible function ϕ and the corresponding Dirichlet integral with the help of our construction with sufficiently large j to arbitrary accuracy.

We see that Courant had little interest in Ritz's method and did not mention Ritz at all. His low esteem for practical questions at that time led Courant even to remove this footnote from the second edition (1925) of the book. He removed, at the same time, the historically first description of what later became one of the most important tools for scientific computations, the finite element method. We will come back to this in Section 7.

Even more sad was the fact that the famous, and by then old, Lord Rayleigh published an article [42], in which he accused Ritz of plagiarism and claimed that all of Ritz's ideas were already present in his own previous work. This led to the name of "Rayleigh-Ritz's" for this method, accepted by many scientists. A very careful study of all the original papers by Leissa [30] made clear that this claim is not at all justified.

6. Immediate use of Ritz's Method in Russia. Contrary to Western Europe, Ritz's method was immediately put to use in Russia, in order to solve hard problems coming from engineering.

6.1. Timoshenko. S.P. Timoshenko (1878–1972), a professor then working at the Politekhnicheskyy Institut of Kiev, later in St. Petersburg and still later at Stanford, was the first to realize the importance of Ritz's invention for applications [44]:

Nous ne nous arrêterons plus sur le côté mathématique de cette question: un ouvrage remarquable du savant suisse, M. Walter Ritz, a été consacré à ce sujet. En ramenant l'intégration des équations à la recherche des intégrales, M. W. Ritz a montré que pour une classe très vaste de problèmes, en augmentant le nombre de paramètres a_1, a_2, a_3, \dots , on arrive à la solution exacte du problème. Pour le cycle de problèmes dont nous nous occuperons dans la suite, il n'existe pas de pareille démonstration, mais l'application de la méthode approximative aux problèmes pour lesquels on possède déjà des solutions exactes, montre que la méthode donne de très bons résultats et pratiquement on n'a pas besoin de chercher plus de deux approximations¹⁹

We also show the 'shveitsarskogo uchenogo Val'tera Rittsa', of this quote in the original Russian publication in Fig. 6.1. Timoshenko then showed how many interesting problems can immediately be solved approximately, using Ritz's method.

¹⁹We will not address the mathematical aspects of this method: a remarkable publication of a Swiss scientist, Mr. Walter Ritz, was dedicated to this subject. Transforming the problem of integrating the equations into a problem of evaluating integrals, Mr. W. Ritz has shown for a large class of problems, that by increasing the parameters a_1, a_2, a_3, \dots , one can find the exact solution of the problem. For the problems we are interested in here however, such a proof does not exist, but the application of the method to problems for which we know an exact solution shows that the method gives very good results and in practice one does not need to compute more than two approximations.



СТРОИТЕЛЬНАЯ МЕХАНИКА КОРАБЛЯ [1]
[РАЗДЕЛЫ, ОТНОСЯЩИЕСЯ К ТЕОРИИ ПЛАСТИН]

Structural Mechanics of Shipbuilding
[Part concerning the theory of shells]

FIG. 6.2. Ivan Bubnov (1872-1919) and the title of his monumental manual for the construction of ships

6.2. Bubnov. Ivan Bubnov (1872-1919, Fig. 6.2), was a structural engineer specialized in the construction of ships, in particular submarines. Like Timoshenko, he was also working at the Polytechnical Institute of St. Petersburg, and needed to calculate the behavior of shells for the construction of submarines. Fascinated by the work of Timoshenko (Bubnov does not cite the work of Ritz directly), the simplicity of the approach and the accuracy of the results, he developed an entire battery of problems with approximate solutions in his manual on ship building [4]. A main contribution of Bubnov to the development of the finite element method is that he realized in [3], after having studied the Zhuranskyi prize-winning work of Timoshenko, that

“... extremely simple solutions can also be obtained in the usual way, i.e., without resorting to a consideration of the energy of the system [...] we simply substitute the expansion for w in the general differential expression for equilibrium, multiply the expression obtained by $\varphi_k dx dy$ and integrate over the entire volume of the body, then we obtain an equation relating the coefficient a_k with all others [...] and will be identical to those found by Prof. Timoshenko.”

This remark simplified the construction of the linear systems for the computation of the coefficients by giving an easy-to-remember recipe. Bubnov also required the coordinate functions to be orthogonal in his remark, without giving a justification. He used in general trigonometric functions for φ_k .

6.3. Galerkin. Boris Grigor’evich Galerkin (1871-1945, Fig. 6.3), was from a poor family, and had to start working already at the age of twelve as a calligrapher. Nevertheless, he managed to study in the Mechanics Department of St. Petersburg Technological Institute, and then worked for the Russian Steam-Locomotive Union and the China Far East Railway. He was also interested in politics, and was arrested in 1905 for political activities, and put into prison for one and a half years. This completely changed his interests, and it was while in prison when he decided to devote the rest of his life to science. Galerkin went on an extensive scientific trip through Europe in 1909, and visited Switzerland, among other countries. We do not know if he met Walter Ritz on this trip, but Galerkin makes precise references to both papers of Ritz in his most famous publication [15], which is now usually quoted when referring to the Galerkin method; Galerkin himself calls this method the Ritz Method, see Fig. 6.4. In this paper, Galerkin introduces Ritz’s method on the first few pages,



СТЕРЖНИ И ПЛАСТИНКИ

РЯДЫ В НЕКОТОРЫХ ВОПРОСАХ УПРУГОГО РАВНОВЕСИЯ
СТЕРЖНЕЙ И ПЛАСТИНОК*

(Петроград, 1915)

Beams and Plates

Using Series for some problems in
the elastic equilibrium of rods and plates

(Petrograd, 1915)

FIG. 6.3. Boris Grigor'evich Galerkin and the famous paper which is now quoted in the literature for the invention of the "Galerkin" method

Из приближенных методов решений широкое применение получил в последнее время метод Ритца^[6, 7]. Этот метод сводится вкратце к следующему.

ЛИТЕРАТУРА

...

3. Бубнов И. Г. Строительная механика корабля. Ч. 2, СПб., 1914.

...

6. Ritz W. Über eine neue Methode zur Lösung gewisser Variationsprobleme der mathematischen Physik.— Journal f. reine u. angewandte Mathematik, 1908, Bd. 135, S. 1—61.

7. Ritz W. Theorie der Transversalschwingungen einer quadratischen Platte mit freien Rändern.— Annalen der Physik, 1909, Bd. 28, N. 4, S. 737—786.

8. Тимошенко С. П. Об устойчивости упругих систем. Киев, 1910, 186 с.

FIG. 6.4. Galerkin himself calls the method which carries today his name the 'Ritz Method', and cites all of Ritz papers, as well as the work by Bubnov and Timoshenko

and then goes on to show how the method can be used to solve approximately many interesting and difficult problems from applications. Galerkin notices that the coordinate functions do not need to be orthogonal, one simply gets an additional matrix, which is today called the mass matrix. However, the main contribution of Galerkin in this paper is to realize that one does not even need a minimization principle in order to construct a finite dimensional system following the recipe given by Bubnov. One can only speculate why the method is nowadays mostly called the Galerkin method; maybe Galerkin's numerous examples from applications were more appreciated than the theoretical analysis of Ritz. Most readers did not even notice the name of the method in Galerkin's introduction, and went straight to the examples, which form the bulk of Galerkin's paper.

7. Late Recognition of Ritz's Work in Western Countries and the Birth of Finite Element Methods. Inspired by Hilbert's ideas, a tremendous number of pure mathematicians applied what became known as direct methods of variational



As Henri Poincaré once remarked, “solution of a mathematical problem” is a phrase of indefinite meaning. Pure mathematicians sometimes are satisfied with showing that the non-existence of a solution implies a logical contradiction, while engineers might consider a numerical result as the only reasonable goal. Such one sided views seem to reflect human limitations rather than objective values.

FIG. 7.1. *Richard Courant and the beginning of his address to the American Mathematical Society in 1941, marking a turning point in his appreciation of Ritz’s work*

calculus in order to prove various existence theorems. An overview of these developments is given in the paper by Richard Courant [10]. In this paper, Courant gives credit to Walther Ritz for the method of finding the coefficients²⁰.

Almost three decades after Ritz’s tragic death, Richard Courant became really interested in practical applications, and gave an address to the American Mathematics Society, on May 3rd, 1941 [11], which he started with the text shown in Fig. 7.1. In this address, Courant gives great praise to the work of Ritz:

“At first, the theoretical interest in existence proofs dominated, and only much later were practical applications envisaged by two physicists, Lord Rayleigh and Walther Ritz. They independently conceived the idea of utilizing this equivalence for numerical calculation of the solutions, by substituting for the variational problems simpler approximating extremum problems in which but a finite number of parameters need be determined”

“But only the spectacular success of Walther Ritz and its tragic circumstances caught the general interest. In two publications of 1908 and 1909, Ritz, conscious of his imminent death from consumption, gave a masterly account of the theory, and at the same time applied his method to the calculation of the nodal lines of vibrating plates, a problem of classical physics that previously had not been satisfactorily treated.”

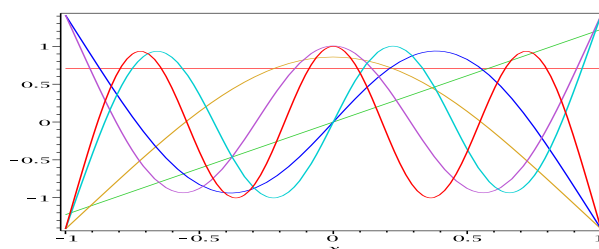
Clearly the tide has turned, and Ritz’s work is now held in high esteem by Courant. In addition, Courant now also realized the importance of the footnote from the first volume of his book with Hurwitz [23]:

“However, the difficulty that challenges the inventive skill of the applied mathematician is to find suitable coordinate functions”.

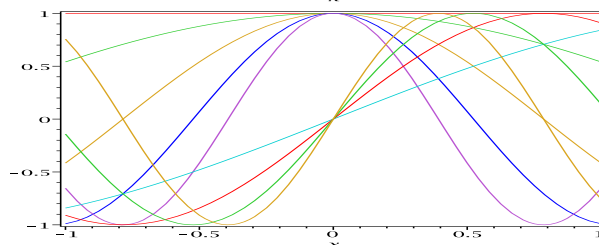
Instead of the eigenfunctions used by Ritz as coordinate functions, and the trigonometric functions employed by Bubnov and Galerkin, Courant now suggests the use of what is nowadays called *hat functions*, see Fig. 7.2. In two dimensions, the piecewise

²⁰ “...einer vielfach auch praktisch angewandten allgemeinen Methode, welche nach Walther Ritz benannt wird.”

Ritz's choice:
Eigenfunctions of the 1d
beam, or polynomials



Bubnov/Galerkin:
Use of trigonometric
functions or polynomials



Courant's choice:
Use hat functions, or
polynomials on elements

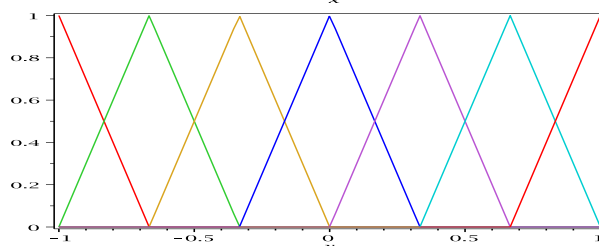


FIG. 7.2. Different choices for the coordinate functions

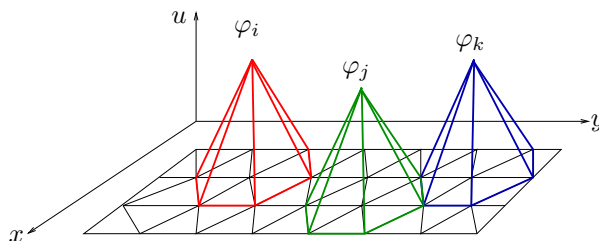


FIG. 7.3. Examples of hat functions in two dimensions

linear functions Courant has in mind are most easily defined on a triangular mesh, in contrast to the up to then dominating finite difference method on a rectangular grid:

“Instead of starting with a quadratic or rectangular net we may consider from the outset any polyhedral surfaces with edges over an *arbitrarily chosen* (preferably triangular) *net*.”

We show a few examples of such hat functions in Fig. 7.3. In his address, Courant presents the first finite element calculation that we were able to find, shown in Fig. 7.4. Courant chooses as a model problem a square domain with a hole, see Fig. 7.4 on the left, and wants to minimize the functional

$$\int \int ((\nabla u)^2 + 2u) dx dy \longrightarrow \min,$$

with $u = 0$ on the outer boundary, and $u = c$, an unknown constant on the inner

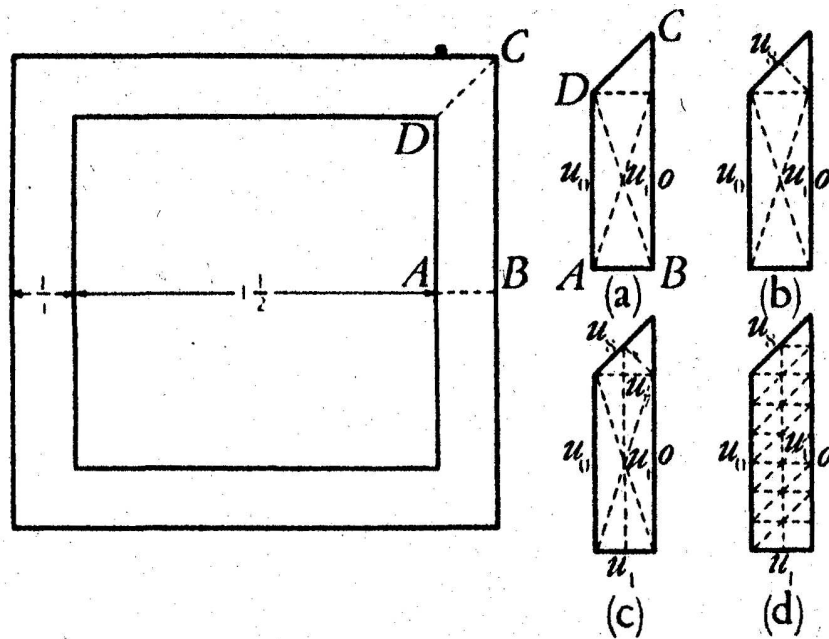


FIG. 7.4. *The first finite element computation we were able to find, by Richard Courant*

boundary. He then compares the results he can obtain for c and the total stiffness

$$S := - \int \int u \, dx \, dy$$

with different coordinate functions:

Polynomials: “with negligible amount of numerical labor”

$$\varphi_1 := a(1 - x) \quad S = 0.339, \quad c = -0.11$$

$$\varphi_2 := a(1 - x)[1 + \alpha(x - \frac{3}{4})y] \quad S = 0.340, \quad c = -0.109$$

Finite Elements: (“The results, easily obtainable, were”)

$$\text{Case (a): } S = 0.344, \quad c = -0.11$$

$$\text{Case (b): } S = 0.352, \quad c = -0.11$$

$$\text{Case (c): } S = 0.353, \quad c = -0.11$$

$$\text{Case (d): } S = 0.353, \quad c = -0.11$$

Courant concludes that “these results show in themselves and by comparison that the generalized method of triangular nets seems to have advantages.”

The development of what was to become the finite element method proceeded on many fronts: on the mathematical side, Trefftz [45] had proposed in 1926 to use local basis functions in the Ritz method. In Russia, Petrov [34] proposed to use approximation spaces different from the testspaces, which led to the now called Petrov-Galerkin family of methods ²¹. A little later, Polya [35] gave an interpretation of the finite difference method using piecewise polynomial approximations, and Argyris [1] combined system theory available at that time with variational approximation in a fundamental step toward a true finite element methodology. More from the engineering side, Hrennikoff [22] used local approximations in the form of bars, beams and spring elements in order to treat plane elasticity problems, and McHenry [32] developed his 'lattice analogy'. During the same time, Kron developed his tensor analysis of networks using a method of tearing, see [27, 28], and Levy developed the direct stiffness method for aircraft simulation [31]. The term *finite element method* was then coined by Ray Clough in [6], who had started to work with Jon Turner from Boeing on structural dynamics, and this work led to the first published description of the finite element method, without the name yet, in [46]; see also the historical note by Ray Clough [7].

8. An Application. We now leave the sunny Swiss Alps from the beginning of this article (Fig. 1.1), to turn to the cold Canadian winter, to illustrate how Ritz's method, and at the same time Schwarz's method, see (2.10), are used today in scientific computing. We compute the temperature distribution in a heated, but less well insulated apartment on Durocher Street in Montreal, shown in Fig. 8.1. The walls are shown in blue, the windows in black on top, and there is also a door in black at the bottom and on the right hand side. We assume that the windows are at $-20^{\circ}C$, and the doors at $15^{\circ}C$. The steady-state heat equation is

$$-\Delta u = f, \quad \text{in the apartment,}$$

where the heat source term f is non-zero at the heater locations, which are close to the window for the two big rooms on the left and the right, on the left wall in the third room from the left, and there is no heater in the second room from the left.

We use a Schwarz domain decomposition method, where the apartment is decomposed into separate rooms, as shown by the red partitions in Fig. 8.1, and a finite element discretization of the problem with linear triangular finite elements, as shown by the mesh in the figure. The "Grundfunktionen" of the Ritz method are therefore simple linear hat functions. In the middle in Fig. 8.1 we show the first iteration of the Schwarz method, where one can clearly see the isolated effect of the heaters and warm doors in each subdomain. At the bottom we show the final result of the simulation, which is now continuous. This result is interesting: one can see that while the heater in the living room on the left and the bedroom on the right is well placed to block the cold from the windows, the heater on the left wall in the bathroom is not effective to keep the room warm, a fact the occupant strongly felt in winter. Also the kitchen is not heated and stays cold, except when cooking and baking.

9. Conclusions. We have tried to give as complete a description as possible of the development which led from the variational calculus to Ritz's work and the

²¹Petrov wrote "Metod Galerkina" for Galerkin's "Metod Ritza". On the other hand, S.L. Sobolev, in the first edition of his textbook *The equations of mathematical physics* from 1947 dedicated an entire chapter to "Metod Ritza". Kantorovich and Krylov in their book *Approximation methods in numerical analysis*, third Russian edition, Moscow 1952, have a section entitled "Metod Ritza i Metod B.G. Galerkina"

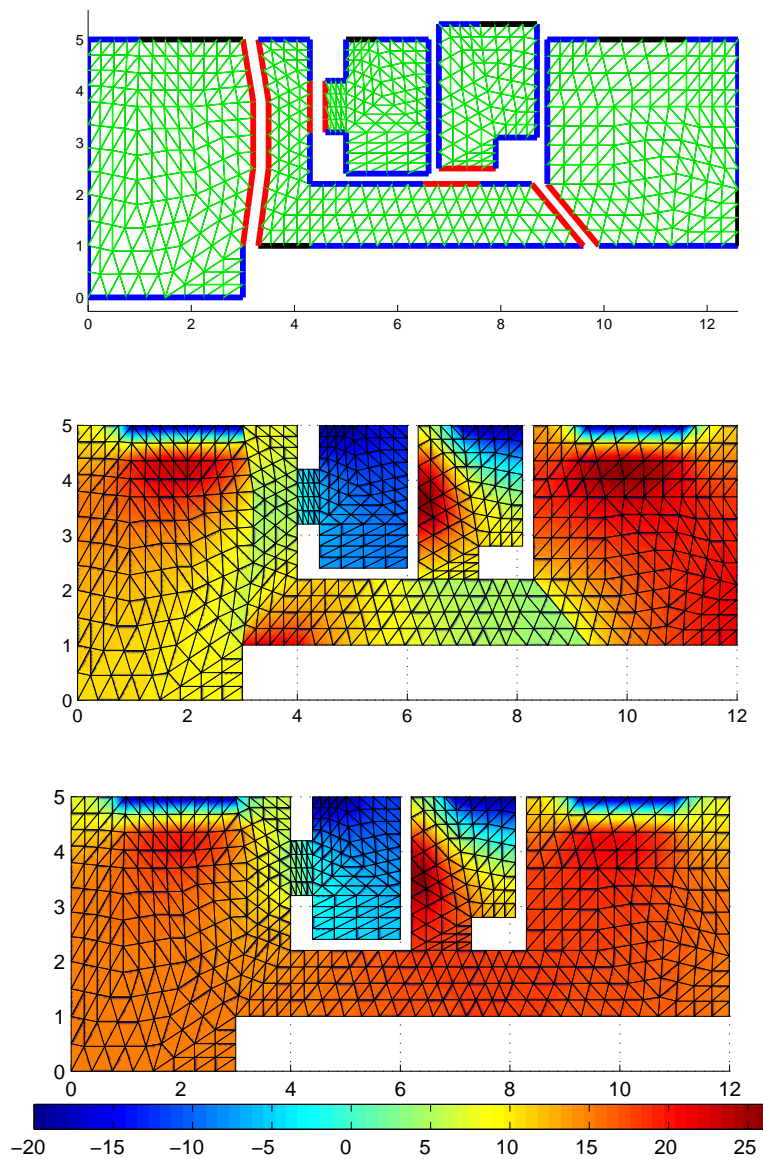


FIG. 8.1. *Triangulation and decomposition of the apartment in Montreal, in the middle the first Schwarz iteration and at the bottom the converged temperature distribution.*

finite element method. There are other descriptions of these historical developments already in the literature, see for example the nice article by Oden [33], which also contains the history of the development of the mathematical theory, the review of the Courant element by Babuška [2], the short description by Taylor [43], or the longer study by Leissa [30]. An extensive 'glimpse' at the literature of variational calculus can be found at the end of volume II of Giaquinta and Hildebrandt [17].

9.1. Back to Euler — the Last Surprise. After all these controversies (should the method be called Rayleigh method? Rayleigh-Ritz? Ritz method? Ritz-Galerkin? Bubnov-Galerkin? Timoshenko-Bubnov-Galerkin? Galerkin method?), let's have a

fresh look at Euler’s original work from 1744: We have heard above that Euler published 1744 the differential equations (1.3) for the variational problem (1.2) and that the standard proof for this equation was found much later. The natural question is thus: *how did Euler himself find “his” differential equation for a variational problem?* We observe in Caput II of E65 the following procedure:

1. Approx. curve by polygon
2. Approx. integral by sum
3. differentiate with resp. to ν
4. set this derivative to zero, apply Euler’s method inversely

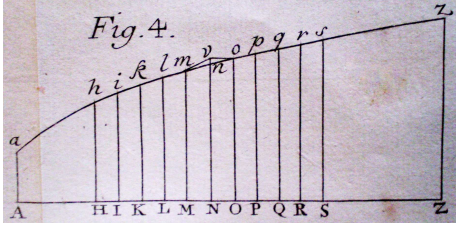


Fig. 4.

$$Z dx + Z' dx + Z'' dx + Z''' dx + \&c.$$

($P + N' dx - P'$) where $N = \frac{\partial Z}{\partial y}$, $P = \frac{\partial Z}{\partial p}$

$$N - \frac{dP}{dx} = 0$$

What did Euler do? In point 1 he discretized the curve by a *finite dimensional* object, which is precisely from the hat function finite element space applied to the 1D case. In point 3 he solved *directly* the finite dimensional problem which he had obtained after discretizing the integral in point 2. Anyway, this *direct* solution by a discretized finite element approximation was the only thing he *could* do, before “his” differential equation in point 4 had been found.

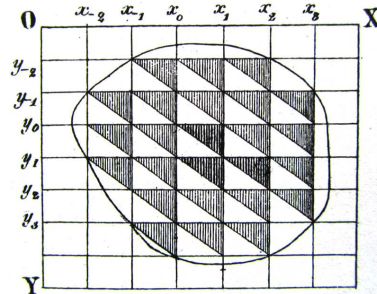


FIG. 9.1. The 2D triangular finite element space in Schellbach 1852.

9.2. Another surprise. Also the 2D triangular finite element approximations can be found in the older literature, in the publication of Schellbach [40] from 1852, who proposed to solve many problems of variational calculus in the spirit of Euler. For his treatment of the minimal surface problem in §30, see Fig. 9.1.

Acknowledgement: the authors would like to thank Jean-Claude Pont, the initiator of the centenary conference of Ritz’ death in Ritz’ hometown Sion, Switzerland. We also are grateful to Ernst Hairer, Ulrich Langer, Fritz Nagel and Heinrich Reitberger for many hints, references and comments, and also to two anonymous referees for reading our manuscript with great care. T. Smirnova, A. Alekseev in Geneva helped us to understand the Russian texts.

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